

# Understanding the White-Emitting $\text{CaMoO}_4$ Co-Doped $\text{Eu}^{3+}$ , $\text{Tb}^{3+}$ , and $\text{Tm}^{3+}$ Phosphor through Experiment and Computation

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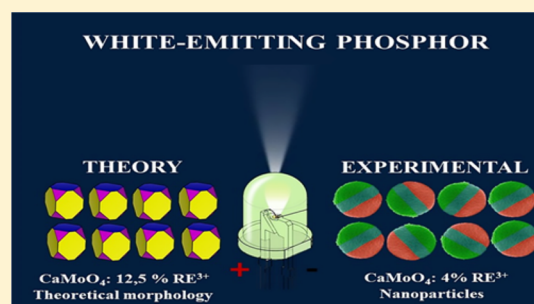
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## Supporting Information

**ABSTRACT:** In this article, the synthesis by means of the spray pyrolysis method, of the  $\text{CaMoO}_4$  and rare-earth cation ( $\text{RE}^{3+}$ )-doped  $\text{CaMoO}_4:x\text{RE}^{3+}$  ( $\text{RE}^{3+} = \text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$ ; and  $x = 1, 2$ , and 4% mol) compounds, is presented. The as-synthesized samples were characterized using X-ray diffraction, Rietveld refinement, field emission scanning electron microscopy (FE-SEM), Raman spectroscopy, and photoluminescence (PL) spectroscopy. To complement and rationalize the experimental results, first-principles calculations, at the density functional theory level, have been performed to analyze the band structure and density of states. In addition, a theoretical method based on the calculations of surface energies and Wulff construction was applied to obtain the morphology transformation of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4:\text{RE}^{3+}$  microstructures. The experimental morphologies can be observed in the FE-SEM images. The PL behavior of the Co-doped samples exhibited well-defined bands in the visible region. The samples with 2 and 4% of  $\text{RE}^{3+}$  released white emission according to the chromaticity coordinates (0.34, 0.34) and (0.34, 0.33), respectively. The present results provide not only a deep understanding of the structure–property relationships of  $\text{CaMoO}_4$ -based phosphor but also can be employed as a guideline for the design of the electronic structure of the materials and the fabrication of photofunctional materials with optimal properties, which allows for the modeling of new phosphors for applications in solid-state lighting.



## 1. INTRODUCTION

Designing new materials for solid-state lighting and understanding the various intricacies involved in designing them such as defects, energy transfer, and concentration quenching is very important. Such materials will be highly beneficial in optical devices, energy, and health industries. The current challenge is to fabricate a material that meets the requirements (thermal and chemical stability, color purity, and emitted color productivity) for its use. White light-emitting diodes (WLEDs) consist of blue LED chips associated with  $\text{YAG}:\text{Ce}^{3+}$  (yellow emission) phosphor,<sup>1</sup> which is considered as an important source in the display technology. However, there are several limiting factors with respect to its use, which correspond to a low color rendering index (CRI) due to the scarcity of red light in the  $\text{YAG}:\text{Ce}^{3+}$  emission, in addition to the low color stability associated with the changes in the transmission voltage. To enhance the CRI, WLEDs are fabricated from a single-phase matrix that efficiently hosts rare-earth type dopants, which exhibit different types of emission in the visible region.<sup>2</sup> Notable characteristics of tridoped matches with rare earth are their luminous efficiency, excellent CRI, adjustable temper-

ature, and purity with respect to chromaticity coordinates.<sup>3</sup> According to Park et al.,<sup>4</sup> the energy transfer (matrix  $\rightarrow$  dopant) coupled with the adjusted concentrations of activators and sensitizers can be worked up to obtain white light-emitting phosphors.

Alkali-earth molybdates  $\text{MMoO}_4$  ( $M = \text{Ca}, \text{Sr}, \text{Ba}$ ) include a wide range of scientifically and technologically significant compounds. These materials are chemically designed and further synthesized using different methods based on the final objective.<sup>5–9</sup> In particular, molybdates phosphors were implemented in optical applications<sup>10,11</sup> and evaluated extensively over the past century as a scintillating medium and electro-optical devices, given their broad and intense absorption bands due to the charge transfer (CT) from oxygen to the metal in the near-ultraviolet (UV) region.<sup>12–16</sup> Among molybdates,  $\text{CaMoO}_4$  is considered an outstanding matrix material for the synthesis of highly applicable luminescent

Received: May 1, 2019

Revised: July 5, 2019

Published: July 9, 2019

materials due to its excellent physical and chemical properties.<sup>17–24</sup> Moreover,  $\text{CaMoO}_4$  exhibits a tetragonal structure and space group  $I4_1/a$ .<sup>25</sup> In the scheelite structure, the Mo cations are bonded to four oxygens, which form the tetrahedral  $[\text{MoO}_4]$  clusters; whereas the Ca cations are coordinated to eight oxygen anions, which form the dodecahedron  $[\text{CaO}_8]$  clusters<sup>10,26</sup> and can be described as a layered arrangement composed of noninterconnected  $[\text{MoO}_4]$  and  $[\text{CaO}_8]$  clusters that form chains parallel to  $[1\bar{1}0]$  with a shared edge.<sup>27</sup> In addition,  $\text{CaMoO}_4$ -based materials are also important materials that can be used in various applications such as photoluminescence (PL) and microwave applications,<sup>28,29</sup> scintillators,<sup>30,31</sup> photocatalysis,<sup>7</sup> and the emission of light in the blue-green region of the electromagnetic spectrum after excitation with UV light.<sup>32</sup> In particular, when cations are introduced into the structure, structural distortions of the tetrahedron and dodecahedron chains occur.<sup>33</sup> Structural disorders in crystals of inorganic materials are present where two or more types of cations or anions statistically occupy the same lattice site.

Inorganic materials doped with  $\text{RE}^{3+}$  ( $\text{RE}^{3+}$ : rare earth) have structural disorders with a distribution of local cationic environments around the doping sites, which deviate from that of the ideal crystallographic positions. These materials attracted significant attention due to their novel structural characteristics and potential applications in various fields.<sup>34,35</sup> In particular,  $\text{RE}^{3+}$  are added to the host matrices to enhance the luminescent properties.<sup>19,36</sup> PL spectroscopy is a very efficient tool for the evaluation of the order–disorder effects in semiconductors, which leads to a better understanding of the structural changes that take place in materials. In particular, the scheelite-type molybdates are interesting materials<sup>37,38</sup> due to their scientific significance, especially with respect to technological applications.<sup>39</sup> The electronic and optical properties of the luminescent materials are significantly influenced by intrinsic and extrinsic defects.<sup>40</sup> These dopant elements favor luminescent properties due to their characteristics and well-defined transitions in the visible range. Moreover, because of their excellent spectroscopic properties, several applications that involve the use of  $\text{RE}^{3+}$  cations have been reported.<sup>33,41–43</sup>

Furthermore,  $\text{CaMoO}_4$  exhibits a broad blue-green luminescence emission peak in the range of 350–650 nm with a peak maximum at approximately 500 nm, which can be tuned as a function of the particle size and the presence of oxygen vacancies in the lattice.<sup>44</sup> In theory, this behavior can be attributed to the electron–hole recombination process after excitation through the band gap of the CT from the 2p orbital of the O anions to the 4d orbitals of the Mo cations in the  $(\text{MoO}_4)^{2-}$  complex tetrahedron.<sup>45</sup> However, the emission band is very broad, and it is difficult to tune the color, particularly for lighting and display applications. On the other hand, the  $\text{RE}^{3+}$  cations have poor absorption cross-sections due to the forbidden nature of the f–f transitions with poor luminescence emission.<sup>46</sup> For the PL emissions, the electron transitions in the 4f and 4d orbitals of the  $\text{RE}^{3+}$  and  $\text{Mo}^{6+}$  cations, respectively, should be considered. Given that the energy level of the 4f orbital is lower than that of the 4d orbital, the possible electron transitions are 4f–4f or 4f–4d. Because of the orbital radius, the 4f orbital is shielded by the 4s and 4p orbitals; however, the 4d orbital is exposed. Therefore, for  $\text{RE}^{3+}$  cations, the 4f–4f transitions are not influenced by the surroundings, and they exhibit a constant energy gap; whereas the 4f–4d transitions are influenced by the structure, and they

exhibit a broad spectrum. The intense and abrupt 4f–4f transitions are a result of the protected external environment due to the 5s and 5p electrons. In addition, according to the Laporte rule, the 4f–4f transition is prohibited; thus, the luminescence based on these transitions is weaker than that of the 4f–4d transitions. Hence, by combining the 4f–4f transitions, which exhibit intrinsic luminescence, with the 4f–4d transitions, which vary depending on the structure, orange-yellow phosphors with a variety of elements can be developed. This leads to an improvement in the light quality and color reproducibility.<sup>47,48</sup> Consequently,  $\text{RE}^{3+}$ -doped materials can be used in various applications such as field emitting displays and LEDs.<sup>49–54</sup>

$\text{CaMoO}_4$  has good chemical and heat stabilities, and it is considered an excellent matrix material. Its luminescence properties can be enhanced when doped with the  $\text{RE}^{3+}$  cations,<sup>19,55–57</sup> e.g.,  $\text{RE}^{3+} = \text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$ . Moreover, structural disorders occur when  $\text{RE}^{3+}$  cations occupy the crystallographic sites of  $\text{CaMoO}_4$ , and when an energy-transfer process takes place from the host/sensitizer to the excited state of the activator by a nonradiative energy-transfer mechanism.<sup>58</sup> In addition, both processes render a high population of photons in the excited level/levels of  $\text{RE}^{3+}$  with concomitant enhanced emissions of  $\text{RE}^{3+}$ .<sup>59</sup> Different synthesis methods have been employed to enrich the luminescent efficiency of the as-prepared luminescent materials, ions-doped molybdates, which include the  $\text{CaMoO}_4$  material.<sup>13,22,32,60–65</sup> The above-mentioned synthesis methods include the Czochralski method,<sup>18</sup> chemical precipitation method,<sup>66,67</sup> citrate–gel method,<sup>68</sup> hydrothermal method,<sup>69–71</sup> polyol process,<sup>72</sup> super-saturated recrystallization process,<sup>73</sup> solid-state reaction,<sup>12</sup> sol–gel process,<sup>13</sup> wherein the quantum efficiency of the material is 20% higher than that obtained using a solid-state reaction sonochemical method,<sup>74</sup> and microwave reaction.<sup>27,75</sup>

The spray pyrolysis method is used for the synthesis of aluminate phosphors and LED phosphor applications, as it offers high efficiency for the formation of non-agglomerated and fine particles with well-defined morphologies, which can increase the brightness and resolution.<sup>76</sup> The spray pyrolysis method is also a promising alternative for the synthesis of rare earth-doped molybdates, as it (i) does not require a high-vacuum environment and (ii) it can be carried out by using simple instrumentation at a low cost.<sup>77</sup> The pyrolysis spray method presents as a synthesis route in which nanoparticles in suspension can be easily obtained and manipulated to realize specific morphologies. What makes aerosol science and technology enabled for areas of high performance in nanotechnology.<sup>78</sup> A major objective of the pyrolysis spray method is to present a single-process step, and the realization of high productivity and high control of the properties of the particles produced.<sup>79</sup> The abovementioned characteristics of the method are of great significance with respect to applications in various industrial fields. The control of the size and morphology of the particles produced is mainly determined by the type of aerosol generator.<sup>80</sup> Consequently, atomization is a major focus of this work. The atomization ultrasound, which was used in the experimental procedure, produces droplets with a fairly uniform size distribution. From recent studies in which the spray pyrolysis technique was used, it was reported that it can be implemented for the production of various powder materials such as inorganic oxides, metals oxides, and nanoparticle oxides, among other compounds.<sup>28,29,81,82</sup>

In this work, an effort has been made in that direction and, for the first time, the synthesis of  $\text{CaMoO}_4$  and  $\text{CaMoO}_4$ -doped rare-earth  $\text{RE}^{3+}$  ( $\text{RE}^{3+} = \text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$ ),  $\text{CaMoO}_4\text{:}x\text{RE}^{3+}$  composites at different percentages,  $x = 1\%$  [ $\text{Eu}^{3+}$  (1%),  $\text{Tb}^{3+}$  (1%), and  $\text{Tm}^{3+}$  (1%)],  $x = 2\%$  [ $\text{Eu}^{3+}$  (2%),  $\text{Tb}^{3+}$  (2%), and  $\text{Tm}^{3+}$  (2%)], and  $x = 4\%$  [ $\text{Eu}^{3+}$  (4%),  $\text{Tb}^{3+}$  (4%), and  $\text{Tm}^{3+}$  (4%)], via the spray pyrolysis method without the addition of any surfactants and templates is reported. The as-synthesized samples were characterized using X-ray diffraction (XRD), field emission scanning electron microscopy (FE-SEM), Raman spectroscopy, and PL spectroscopy. The color coordinates, color temperature, and CRI were also investigated. The materials obtained were tested for their applicability as white light emitters. In addition, first-principles quantum-mechanical calculations based on the density functional theory (DFT) are performed to complement and rationalize the experimental results, in particular, the local structure, morphology, band gap, and electronic and optical properties of the pure  $\text{CaMoO}_4$  and  $\text{CaMoO}_4\text{:RE}^{3+}$  composites. A theoretical method, based on the calculations of surface energies and Wulff construction, was applied to study the morphology transformations. From an analysis of the results of the characterization experiment and first-principles calculations, the nature of the mechanism for the luminescence emissions were revealed.

## 2. EXPERIMENTAL SECTION

**2.1. Materials.** Molybdc acid ( $\text{H}_2\text{MoO}_4$ , 85%, Alfa Aesar), calcium nitrate ( $\text{Ca}(\text{NO}_3)_2 \cdot 4\text{H}_2\text{O}$ , 99%, Alfa Aesar), nitric acid ( $\text{HNO}_3$ -synth, 99.9%), europium oxide ( $\text{Eu}_2\text{O}_3$ , 99%, Aldrich), terbium nitrate pentahydrate ( $\text{Tb}(\text{NO}_3)_5 \cdot 5\text{H}_2\text{O}$ , 99% Aldrich 99.9%), thulium oxide ( $\text{Tm}_2\text{O}_3$ , 99%, Aldrich), and distilled water were used as-received for the preparation of the  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructures.

**2.2. Preparation of  $\text{CaMoO}_4\text{:RE}^{3+}$ .** Calcium molybdate powders ( $\text{CaMoO}_4$ ) were processed using the spray pyrolysis method with the addition of europium ( $\text{Eu}^{3+}$ ), thulium ( $\text{Tm}^{3+}$ ), and terbium ( $\text{Tb}^{3+}$ ) rare earths. The experimental procedure was conducted as described. The precursors of calcium nitrate and molybdc acid were dissolved separately in a beaker that contained 100 mL of deionized water. The resulting solution was then added to a beaker, and then homogenized at approximately 60 °C. Under constant stirring and a stabilized temperature (60 °C),  $\text{Eu}^{3+}$  (europium),  $\text{Tb}^{3+}$  (terbium), and  $\text{Tm}^{3+}$  (thulium) were added to the reaction in the proportions presented in Table 1. For the  $\text{Eu}^{3+}$  and  $\text{Tm}^{3+}$

**Table 1.**  $\text{CaMoO}_4\text{:RE}^{3+}$  Compositions Investigated

samples	$\text{RE}^{3+}$		
	$\text{Eu}^{3+}$ (%)	$\text{Tb}^{3+}$ (%)	$\text{Tm}^{3+}$ (%)
$\text{CaMoO}_4$	0	0	0
$\text{CaMoO}_4\text{:}1\%(\text{RE}^{3+})$	0.34	0.33	0.33
$\text{CaMoO}_4\text{:}2\%(\text{RE}^{3+})$	0.67	0.67	0.66
$\text{CaMoO}_4\text{:}4\%(\text{RE}^{3+})$	1.34	1.33	1.33

dopants, it was necessary to dissolve the oxides  $\text{Tm}_2\text{O}_3$  and  $\text{Eu}_2\text{O}_3$  in approximately 10 mL of  $\text{HNO}_3$  under stirring at 60 °C to form the respective nitrates. In this manner, all dopants in the reaction could be solubilized. The reagents were stoichiometrically weighed and added to the calcium molybdate solution. The solution remained under agitation and was temperature-controlled for approximately 30 min.

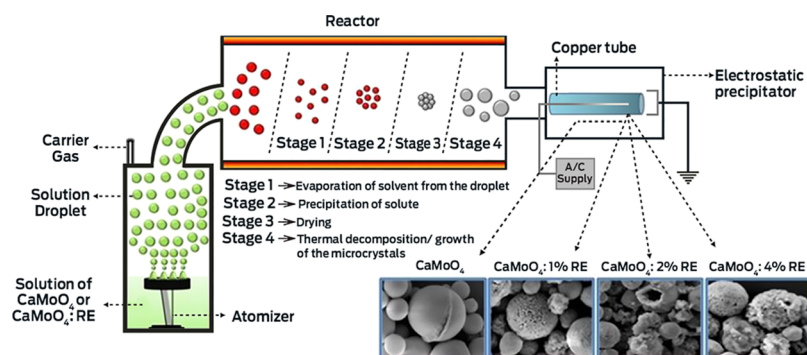
Figure 1 presents a schematic of the spray pyrolysis system. The droplet generator comprised an ultrasonic nebulizer. The laminar flow aerosol reactor used in the present study was a quartz tube with an inner diameter of 37 mm and length of 1.86 m, which was inserted into a horizontal electric furnace. The precursor solution was atomized at a frequency of 2.4 MHz using the ultrasonic nebulizer. The sprayed droplets were carried to the reactor, heated by an electric furnace at 600 °C in air, and converted into solid oxide particles within the laminar-flow aerosol reactor. The resulting particles were collected at the reactor exit using an electrostatic precipitator. The overall flow rate of the air used as a carrier gas was 2 L  $\text{min}^{-1}$ . The residence time of the  $\text{CaMoO}_4\text{:RE}^{3+}$  solution in the reactor is a function of the drag gas flow and the reactor area volume according to eq 1.<sup>83</sup>

$$T = \frac{R}{60Q} \quad (1)$$

where  $T$  is the residence time (s),  $R$  is the gas flow rate (L  $\text{min}^{-1}$ ), and  $Q$  is the reactor volume (L). The reactor volume was 1.61 L, and the air flow rate was 2 L  $\text{min}^{-1}$ . From this relationship, the mean time of the synthesis was calculated as approximately 48 s.

Figure 1 presents a schematic that includes all stages of the synthesis and growth of the nanocrystal  $\text{CaMoO}_4$  and  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructures obtained using the spray pyrolysis method. At Stage 1, the thermal decomposition occurs directly at the nucleation sites, which promotes the crystallization kinetics of the initial crystals. Because of the free rotation of the primary crystals, random collisions and rearrangements of the crystals can occur through movement at this stage, which comprises Stages 2 and 3, and the nanocrystals are effectively generated. At Stage 4, the nanocrystals tend to undergo self-assembly, possibly via van der Waals interactions,<sup>46</sup> to achieve the minimum energy per volume ratio.<sup>84</sup> Favoring the growth kinetics of nanocrystals, the thermal decomposition can promote the heterogeneous nucleation and aggregation of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4\text{:RE}^{3+}$  nanocrystals. The growth promotion occurs at Stage 5, which is responsible for the formation and growth of the aggregated  $\text{CaMoO}_4$  or  $\text{CaMoO}_4\text{:RE}^{3+}$  nanocrystals. The high drying temperature can generate a temperature gradient between the inner surface of the crystals and their interior, which leads to thermal stress on the crystals.

**2.3. Characterization of the  $\text{CaMoO}_4\text{:RE}^{3+}$  Microstructures.** The  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructures were structurally characterized by XRD using a Shimadzu XRD 7000 instrument with Cu  $K\alpha$  radiation ( $\lambda = 1.5406 \text{ \AA}$ ) at  $2\theta = 10\text{--}80^\circ$  and a scanning rate of  $0.02^\circ \text{ s}^{-1}$ . Micro-Raman spectroscopy measurements were carried out using a T-64000 spectrometer (Jobin Yvon, France) with a triple monochromator coupled to a charge couple detector. The spectra were obtained using an argon ion laser with a wavelength of 514.5 nm and a maximum output power of 8 mW. The morphologies were evaluated using field emission-scanning electron microscopy (FE-SEM) (Carl Zeiss, Supra 35-VP Model, Germany). The UV–vis diffuse reflectance spectrum was measured at room temperature using a UV–vis spectrometer. The PL spectra were obtained by using an Ash MonoSpec 27 monochromator (Thermal Jarrel, U.S.A.) and a R4446 photomultiplier (Hamamatsu Photonics, U.S.A.). The 350 nm beam of a krypton ion laser (Coherent Innova 90 K) was used as the excitation source with a maximum output



**Figure 1.** Schematic diagram of the experimental apparatus and the formation of  $\text{CaMoO}_4$  and  $\text{CaMoO}_4:\text{RE}^{3+}$  microstructures obtained by one-pot ultrasonic spray pyrolysis.

power of 200 mW. All measurements were carried out at room temperature.

In the synthesis of  $\text{CaMoO}_4$  and  $\text{CaMoO}_4:\text{RE}^{3+}$ , the quantities of the reagents used were weighed stoichiometrically according to the values presented in Table 1. To verify the chemical composition of the samples, a chemical analysis of the material was carried out from energy-dispersive X-ray spectroscopy (EDS).

### 3. COMPUTATIONAL DETAILS

The bulk of the  $\text{CaMoO}_4$  structure was doped with 12.5% mol of  $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$ . Simulations were conducted using quantum-mechanical calculations based on the DFT framework at the B3LYP hybrid functional level<sup>85,86</sup> in the CRYSTAL17 computer code.<sup>87</sup> This method was successfully employed, as reported in several studies conducted on the bulk, electronic, and structural properties of molybdate-, tungstate-,<sup>91,92</sup> vanadate-<sup>93,94</sup> and perovskite<sup>95</sup>-based materials, in addition to  $\text{Ag}_2\text{CrO}_4$ .<sup>96</sup>

The atomic centers were described using standard all-electron basis sets (6-31G\* basis set) for the O and Ca atoms; whereas the Mo, Eu, Tb, and Tm atoms were described by pseudopotential basis sets.<sup>97</sup> According to the *f*-in-core approximation, the electrons of the 4f shells of  $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$  were incorporated in the pseudopotential. Thus, an explicit treatment of the open 4f shell was not required, which is a major advantage from the computational viewpoint. In the density matrix diagonalization of the bulk calculations, the reciprocal space net was described by a shrinking factor of 4, which corresponds to 36 *k*-points generated according to the Monkhorst–Pack scheme.<sup>98</sup> The accuracy of the evaluation of the Coulomb and exchange series was controlled by five thresholds with values of  $10^{-6}$ ,  $10^{-6}$ ,  $10^{-6}$ ,  $10^{-6}$ , and  $10^{-12}$ .

The initial cell and atomic position parameters used in the optimization process were obtained from the results of the Rietveld refinement of  $\text{CaMoO}_4$  using a conventional unit cell that contained 24 atoms.<sup>10</sup> The calculation of the equilibrium geometries and electronic properties of the  $\text{CaMoO}_4:\text{RE}^{3+}$  (12.5%-doped) system was used to simulate the tetragonal supercell structure of 95 atoms, which corresponds to  $2 \times 2 \times 2$  conventional cells, in which two  $\text{Ca}^{2+}$  ions were replaced by two  $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$  cations, and a calcium vacancy ( $V_{\text{Ca}}$ ) was created to neutralize the unit cell. It should be noted that for the calculation of the doped materials ( $\text{CaMoO}_4:\text{RE}^{3+}$ ), the experimental percentages of doping were 1, 2, and 4 mol %, as it was necessary to use very large unit cells. The calculations were therefore computationally complex, and a minimum

doping amount of 12.5% could be achieved separately for the  $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$  cations.

The equilibrium shape of the  $\text{CaMoO}_4:\text{RE}^{3+}$  microstructures could be determined using the classic Wulff construction,<sup>99</sup> which minimizes the total surface free energy at a fixed volume, and provides a simple relationship between the surface energy,  $E_{\text{surf}}$  of the (*hkl*) plane and its distance in the normal direction from the center of the crystallite. After all atoms were fully relaxed and the system reached the minimized energy, the surface energies,  $E_{\text{surf}}$  could be calculated using the following expression.<sup>92,100</sup>

$$E_{\text{surf}} = (E_{\text{slab}} - E_{\text{bulk}})/2A \quad (2)$$

where  $E_{\text{slab}}$  and  $E_{\text{bulk}}$  are the total energies of the slab and the bulk with the same number of atoms, respectively, and  $A$  represents the surface area of the slab model. The strong dependence of  $E_{\text{surf}}$  on the surface orientation can be understood with respect to the different distributions of the surface atoms on different surface orientations, given that atomic distribution has a significant influence on the surface electronic structure.

The Wulff construction was successfully used in materials science to obtain the morphologies of materials such as  $\text{BaMoO}_4$ ,<sup>90</sup>  $\text{PbMoO}_4$ ,<sup>101</sup> and  $\beta\text{-Ag}_2\text{MoO}_4$ .<sup>102</sup> To confirm the convergence of the total energy with respect to the slab thicknesses of different surface models, the  $E_{\text{surf}}$  values for several low-index surfaces were calculated using a  $2 \times 2$  supercell from pure surfaces, and by substituting two  $\text{Ca}^{2+}$  ions with two  $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$  or  $\text{Tm}^{3+}$  cations, thus creating a  $\text{Ca}^{2+}$  vacancy, as expressed in the bulk calculations.

### 4. RESULTS

The XRD patterns of the samples are presented in Figure 2. All of the diffraction patterns detected were indexed as scheelite-type tetragonal crystal structures with unit cells in the space group  $I4_1/a$  with  $C_{4h}$  symmetries, which is in accordance with the crystallographic card JCPDS 60-5552. This indicates that the rare-earth clusters introduced into the matrix did not induce any significant changes in the  $\text{CaMoO}_4$  structure; however, the density of negative electronic defects was significantly increased. The  $\text{RE}^{3+}$  ion, when incorporated in the  $\text{CaMoO}_4$  matrix, results in an excess of positive charges. Hence, the occurrence of  $\text{Ca}^{2+}$  vacancies is for the neutralization of the charges. As shown in Figure 2, there was a shift in the peak of (112) to the low angle region due to the increase in the concentration of the  $\text{RE}^{3+}$  ions. This behavior may be verified by the difference between the

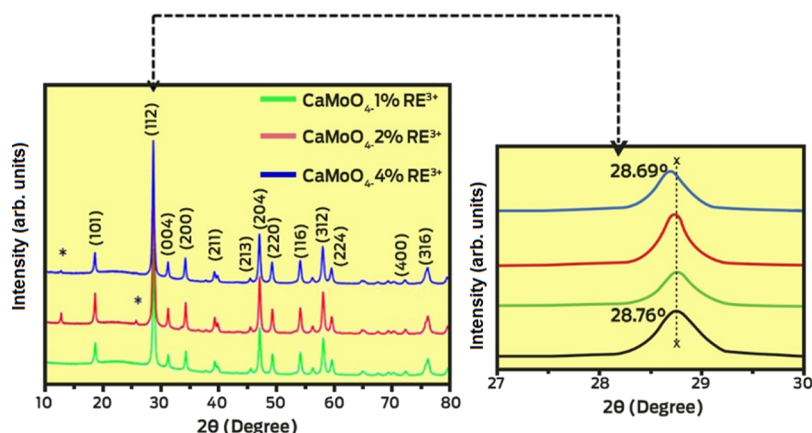


Figure 2. XRD patterns  $\text{CaMoO}_4$  and  $\text{CaMoO}_4:\text{RE}^{3+}$  ( $x = 1, 2,$  and  $4$  mol %) obtained by the spray pyrolysis method.

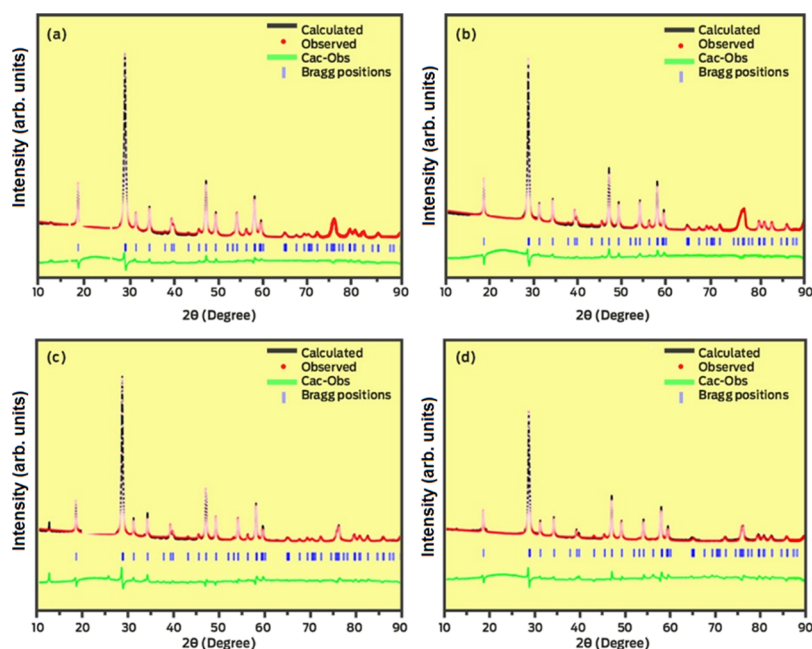


Figure 3. Rietveld refinements of (a)  $\text{CaMoO}_4$ , (b)  $\text{CaMoO}_4:1\%(\text{RE}^{3+})$ , (c)  $\text{CaMoO}_4:2\%(\text{RE}^{3+})$ , and (d)  $\text{CaMoO}_4:4\%(\text{RE}^{3+})$  microstructures.

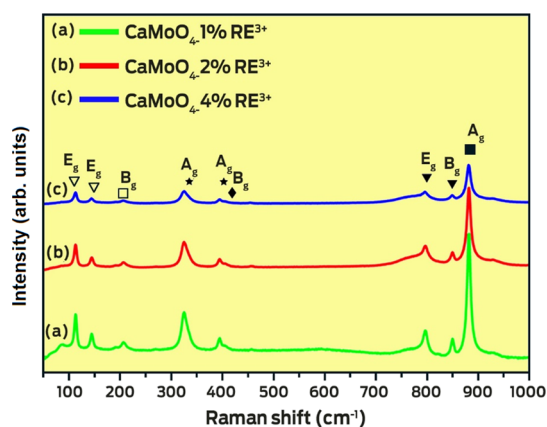
electronic density of the  $\text{Ca}^{2+}$  in relation to the  $\text{RE}^{3+}$ , which induces distortions and therefore polarization in the clusters  $[\text{CaO}_8]$ . These changes occur due to the difference in the electronic densities between the forming clusters of the lattice and the doping clusters. The doping process induces the formation of  $V_{\text{Ca}}$  negative sites that modify the position of the valence band (VB) energy levels, with concomitant insertion of 4f orbitals of  $\text{RE}^{3+}$  cations into the conduction band minimum. As proposed by Almeida et al.<sup>103</sup> and Parchur and Ningthoujam<sup>42</sup> from the Rietveld refinement conducted on the  $\text{CaEu}_2(\text{WO}_4)_4$  and  $\text{CaEu}_2(\text{MoO}_4)_4$  red phosphors, the  $[\text{MoO}_4]^{2-}$  tetrahedra in scheelite-structures have high flexibilities. The metal–oxygen distances and bond angles vary significantly with changes in the population of the A-site (in this work A = Ca) by cations with different charges/sizes and/or cation vacancies.<sup>104</sup> However, as shown in Figure 10, no major long-term structural changes were observed. Additional peaks identified at  $2\theta = 12.87$  and  $36.65$  in the samples doped at 2 and 4 mol % of  $\text{RE}^{3+}$ , respectively, were observed. These peaks can be attributed to the compound  $\text{MoO}_n \cdot m\text{H}_2\text{O}$  formed because of the excessively high solubility

of the  $\text{RE}^{3+}$  in the  $\text{CaMoO}_4$  matrix, which exceeded the limit.<sup>42</sup> The same tendency was observed in the formation of the secondary phase when 7 and 10 mol % of  $\text{Eu}^{3+}$  were added to the  $\text{CaMoO}_4$  structure by the hydrolysis method using urea. The Rietveld refinement method was employed to verify possible differences in the structural arrangements obtained using Maud version 2.0 software, and the results of the analysis are presented in Figure 3.

The measured diffraction patterns were in good agreement with the JCPDS 60-5552 standard. An analysis of the diffractograms of Figure 3 and the values of the refinement parameters in Table S1 (Supporting Information) confirmed the match between the experimental and theoretical data. The results of the refinements are summarized in Table S1. In addition, the close correspondence between the lattice parameters and the unit cell volumes reported in the literature<sup>98</sup> indicate that  $\text{CaMoO}_4$  powders have a scheelite-type structure (tetragonal). The introduction of a dopant into a matrix can promote changes in the intensity and positions of the diffraction peaks, as shown in Figure 2. These changes

occur because of the difference in size between the forming atom of the lattice and the dopant.

Moreover, the scheelite-type tetragonal structure in the primitive cell exhibits 26 different vibration modes:  $\Gamma_{Td} = {}^3A_g + {}^5A_u + {}^5B_g + {}^3B_u + {}^5E_g + {}^5E_u$ . However, only  $A_g$ ,  $B_g$ , and  $E_g$  are Raman-active, whereas the odd modes  ${}^4A_u$  and  ${}^4E_u$  can only be detected in the infrared spectra. The three  $B_u$  vibrations are silent modes. In addition, one  $A_u$  mode and one  $E_u$  mode are acoustic vibrations. The micro-Raman spectra of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructures are presented in Figure 4.



**Figure 4.** Raman spectra for  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructure materials prepared by the pyrolysis spray method.

The primitive  $\text{CaMoO}_4$  cell includes two formula units with a weak coupling between the  $[\text{MoO}_4]^{2-}$  moiety and the  $\text{Ca}^{2+}$  cations, in addition to the presence of strong covalent Mo–O bonds. The external or lattice phonons correspond to the motion of the  $\text{Ca}^{2+}$  cations and the rigid molecular unit. The  $[\text{MoO}_4]^{2-}$  tetrahedral ion in the free space exhibits a  $T_d$  symmetry. The internal vibrational ions correspond to the vibrations within the  $[\text{MoO}_4]^{2-}$  group, with an immovable center of mass.<sup>105</sup>

The Raman spectra revealed well-defined external and internal peaks for the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructures, which indicate that the synthesized powders were highly crystallized. The modes of vibration of the scheelite phase in the tetrahedral structure were observed for each sample. Moreover, several peaks were observed, which were associated with the Raman-active internal modes of tetrahedral  $\text{MoO}_4$ :  $\nu_1$  ( $A_g$ ),  $\nu_3$  ( $B_g$ ),  $\nu_3$  ( $E_g$ ),  $\nu_4$  ( $E_g$ ),  $\nu_4$  ( $B_g$ ),  $\nu_2$  ( $B_g$ ),  $\nu_2$  ( $A_g$ ), R ( $A_g$ ), R ( $E_g$ ), and external T ( $B_gE_gE_g$ ). This is in accordance with the findings reported in refs 20 and 106. Several differences between the Raman spectra of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructures were observed, and these differences can be associated to the changes of the Mo–O bond distances provoked by the addition of  $\text{RE}^{3+}$  to the structure of the  $\text{CaMoO}_4$  lattice.

In Figure 5a–c the unit cell and cluster coordination of the scheelite-type  $\text{CaMoO}_4\text{:RE}^{3+}$  (12.5%-doped) employed in this work have been depicted.

An analysis of the results revealed that the substitution of  $\text{Ca}^{2+}$  by  $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , or  $\text{Tm}^{3+}$  provokes changes of the atomic coordinates of the oxygen anions with concomitant variations in the  $\text{RE}^{3+}$ –O distances (see Figure 5a–c), which indicates the existence of structural and electronic distortions in the  $[\text{REO}_8]$ ,  $[\text{MoO}_4]$ , and  $[\text{CaO}_8]$  clusters. The corresponding values of atomic coordinates ( $x$ ,  $y$ ,  $z$ ) are listed in Table S2.

The oxygen anions (O1–O6), highlighted in Figure 5a–c, allow to find a relationship between the bond distances of the oxygen atoms and the  $\text{RE}^{3+}$  cations in the  $\text{CaMoO}_4$  structure. It should be noted that the presence of different  $\text{RE}^{3+}$  cations produces slight structural distortions in the  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructures, which shortens the  $\text{Eu}^{3+}$ –O and  $\text{Tm}^{3+}$ –O bond distance, as exemplified by the following distance: O1–Eu1 (2.45 Å) to O1–Tm1 (2.37 Å).

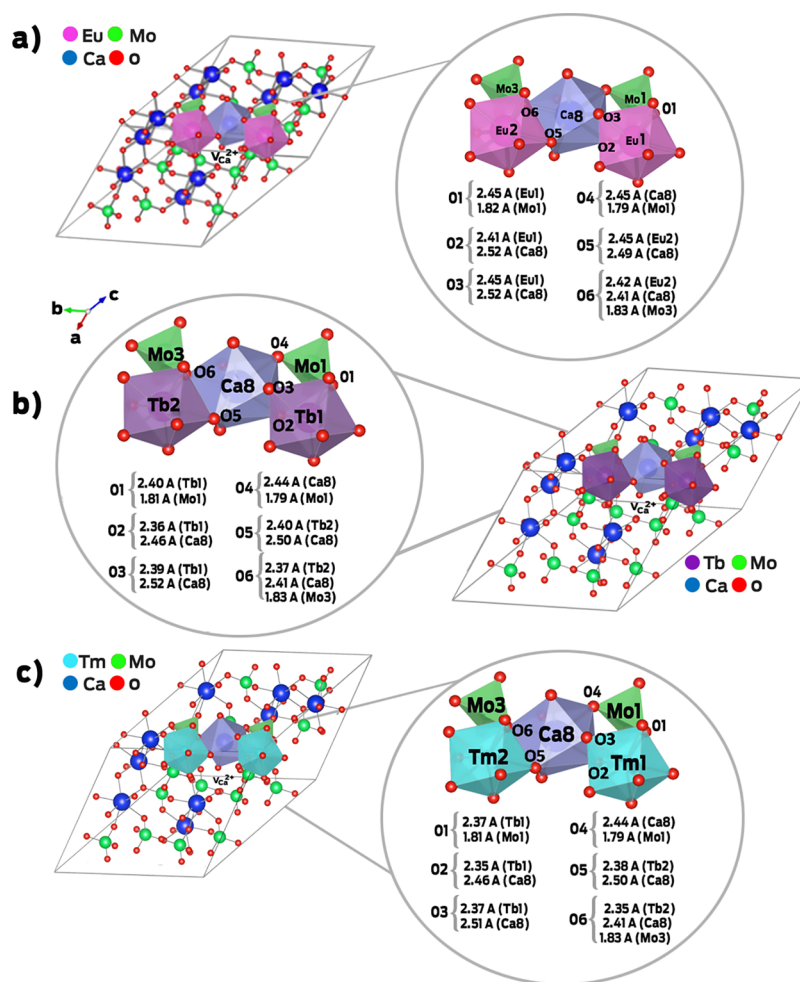
The calculated band structure and density of states (DOS) projected for the atoms of  $\text{CaMoO}_4\text{:RE}^{3+}$  (12.5%) are displayed in Figure 6a–c and the first Brillouin zone of scheelite is shown in Figure S1. The projected DOS for  $\text{RE}^{3+} = \text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$  (Figure 6a–c right panel) reveal that the upper levels of the VB mainly consist of O 2p orbitals; whereas the conduction band (CB) is predominantly formed by  $\text{Mo}^{6+}$  4d and  $\text{RE}^{3+}$  4f orbitals, with a small amount of  $\text{Ca}^{2+}$  orbitals. In addition, the doping was found to cause a decrease in the energy levels of the CB with respect to the fundamental band gap of pure  $\text{CaMoO}_4$  ( $\sim 4.9$  eV);<sup>10</sup> the calculated value is similar to that obtained by Panchal et al.<sup>107</sup> and Ryu et al.<sup>108</sup> The results reveal that both the VB maximum and CB maximum are sensitive to the substitution of the  $\text{RE}^{3+}$  cation electronic states, thus creating intermediate energy levels that can be associated to the presence of 4f orbitals at the CB maximum. In particular, intermediate electronic levels were introduced into the  $E_{\text{gap}}$  region, thus reducing the energy value, as shown in the left panel of Figure 6b,c. An indirect transition was produced along the  $k$ -points  $Z(111)$  to  $\Gamma(000)$  from the top of the VB to the bottom of the CB in  $\text{CaMoO}_4\text{:Tb}^{3+}$  ( $E_{\text{gap}} = 1.17$  eV), and an indirect transition was observed along the  $k$ -points  $X(001)$  to  $\Gamma$  in  $\text{CaMoO}_4\text{:Tm}^{3+}$  ( $E_{\text{gap}} = 3.97$  eV), although the intermediate level in the band structure for the doped system was flat. However, in the case of  $\text{CaMoO}_4\text{:Eu}^{3+}$  ( $E_{\text{gap}} = 4.69$  eV), a direct transition was produced along the  $k$ -points  $\Gamma$  to  $\Gamma$  from the top of VB to the bottom of the CB (see Figure S1).

The experimental values obtained for the fundamental band gap energies of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4\text{:RE}^{3+}$  microstructures were estimated from the respective diffuse-reflectance spectra by plotting the square of the Kubelka–Munk function (i.e.,  $F(R)^2$ ) as a function of the energy (in eV).<sup>109,110</sup> The values were determined by extrapolating the linear part of the curve to  $F(R)^2 = 0$ , as shown in Figure 7. The ratio between the molar absorption coefficient ( $k$ ) and scattering coefficient ( $s$ ) was estimated from reflectance data using the Kubelka–Munk relationship, as expressed by eq 3

$$F = \frac{K}{S} = \frac{(1 - R)^2}{2R} \quad (3)$$

where  $R$  is the percentage of reflected light. The incident photon energy ( $h\nu$ ) and optical band gap energy ( $E_g$ ) are related to the transformed Kubelka–Munk function  $[F(R)h\nu]^n = A(h\nu - E_g)$ , where  $E_g$  is the fundamental band gap energy,  $A$  is a constant dependent on the transition probability, and  $n$  is the power index related to the optical absorption process. Moreover,  $n = 1/2$  or  $n = 2$  for an indirect or direct allowed transition,<sup>110,111</sup> respectively. The corresponding transition for  $\text{CaMoO}_4$  was considered as direct. The  $E_g$  values are presented in Figure 7.

The UV–vis spectra are presented in Figure 7, and the value of the band gap energy was within the range of 3.84–3.93 eV. The substitution of  $\text{Ca}^{2+}$  by  $\text{RE}^{3+}$  leads to an excess of positive charges in the semiconductor structure. To neutralize these



**Figure 5.** Theoretical representation of the scheelite type structure tetragonal  $2 \times 2 \times 2$  supercell corresponding to (a)  $\text{CaMoO}_4$ :12.5% mol  $\text{Eu}^{3+}$ , (b)  $\text{CaMoO}_4$ :12.5% mol  $\text{Tb}^{3+}$ , and (c)  $\text{CaMoO}_4$ :12.5% mol  $\text{Tm}^{3+}$  microstructures.

charges, defects of the  $\text{Ca}^{2+}$  vacancy type (p) occur, and the presence of these defects is responsible for the intermediate levels in the band gap. These levels favor the electronic transitions between the VB and the CB due to the band gap reduction.

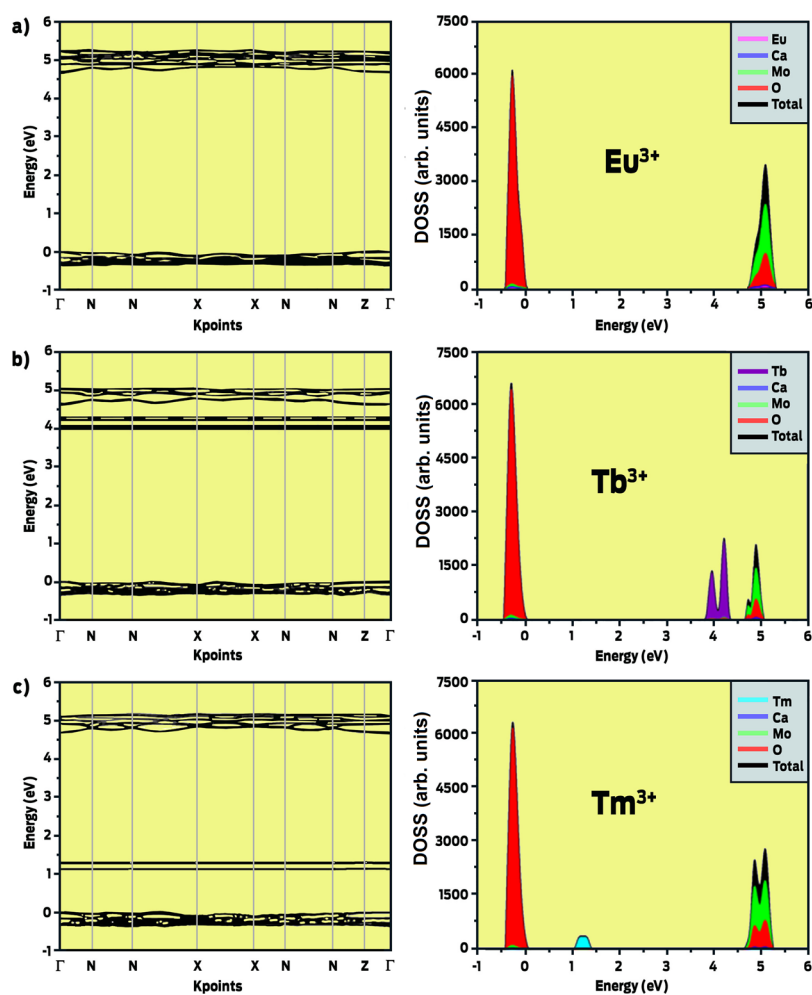
The PL spectra of the pure  $\text{CaMoO}_4$  and  $\text{CaMoO}_4$ : $\text{RE}^{3+}$  are presented in Figure 8a–c, respectively. It is well known that PL emissions<sup>112</sup> can be associated to the structural and electronic distortion in the tetrahedral  $[\text{MoO}_4]^{-2}$  cluster<sup>66</sup> with concomitant appearance of intermediate levels of the band gap.<sup>113</sup> For  $\text{CaMoO}_4$  microstructures obtained at 600 °C, a significant number of structural defects in the crystalline lattice is expected, which is considered a critical factor for the PL properties. Materials with lower crystallinity indexes exhibit a significant number of intermediate energy levels in the band gap region.<sup>114</sup> Under this condition, electronic transitions occur more easily, thus favoring the PL emissions.

An analysis of the results of Figure 8a shows that the PL spectra are composed of broad-band emissions covering the visible electromagnetic spectrum in the range 400–650 nm. This kind of emission profile is typical of multilevel and multiphonon processes, where several paths involving the participation of multiple energy states exist within the band gap. The emission spectra were deconvoluted in order to further elucidate the emission centers and to qualitatively determine the contribution of each visible light component.

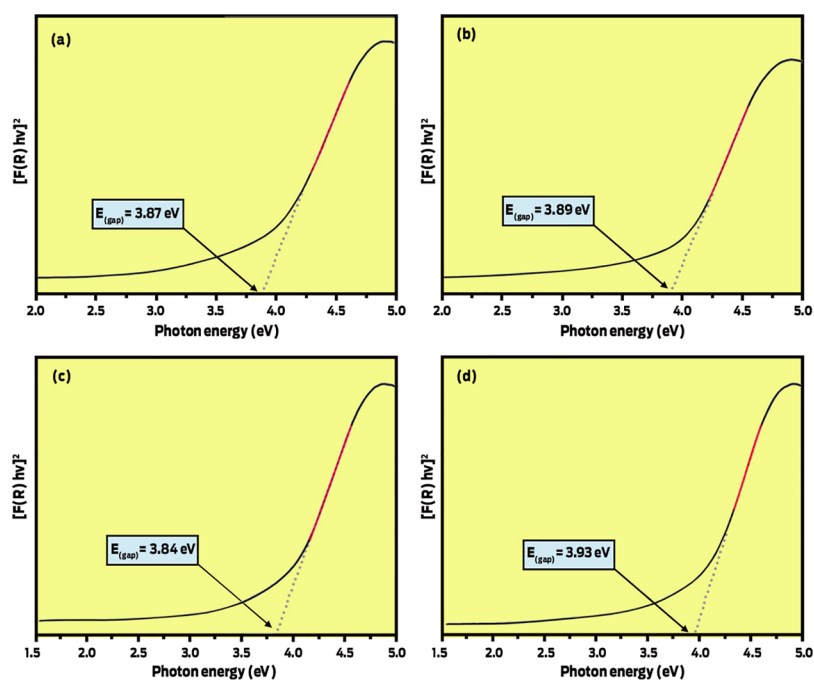
Figure 8b illustrates the results obtained from the deconvolution of the PL profiles.<sup>73,89,103</sup> The profile is well-described by the Voigt area function, and an analysis of the results revealed that the percentages of the areas in the green, blue, orange, and red region are approximately 55, 20, 19, and 6%, respectively.

The presence of the dopant ions ( $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$ ) in the  $\text{CaMoO}_4$  matrix were observed. By mapping the distribution of the elements, it is possible to verify the homogeneous distribution of the rare-earth ions ( $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$  and  $\text{Tm}^{3+}$ ) in all particles of  $\text{CaMoO}_4$ . This homogeneity in the material is of great importance for the efficiency of the PL behavior. In addition, the EDS spectrum and the mapping of the distribution of the atoms has been included (see Figure S3, Supporting Information).

The PL phenomenon is reported to be significantly dependent on the presence of defects in the lattice, which typically arise during the synthesis or heat treatment of the material. Moreover, the  $\text{CaMoO}_4$ : $\text{RE}^{3+}$  material with scheelite-type structure demonstrated a strong  $4f-4f$  transitional absorption and emission due to the polarization of the  $[\text{MoO}_4]^{2-}$  moiety.<sup>42,115</sup> As discussed in DOS and band structure, the  $\text{RE}^{3+}$  doping process at  $\text{CaMoO}_4$  material promotes a decrease in the band gap values, as well as enhance the CT process from the tetragonal  $[\text{MoO}_4]^{2-}$  moieties to the  $[\text{REO}_8]$  clusters, as other authors have pointed out.<sup>116</sup>

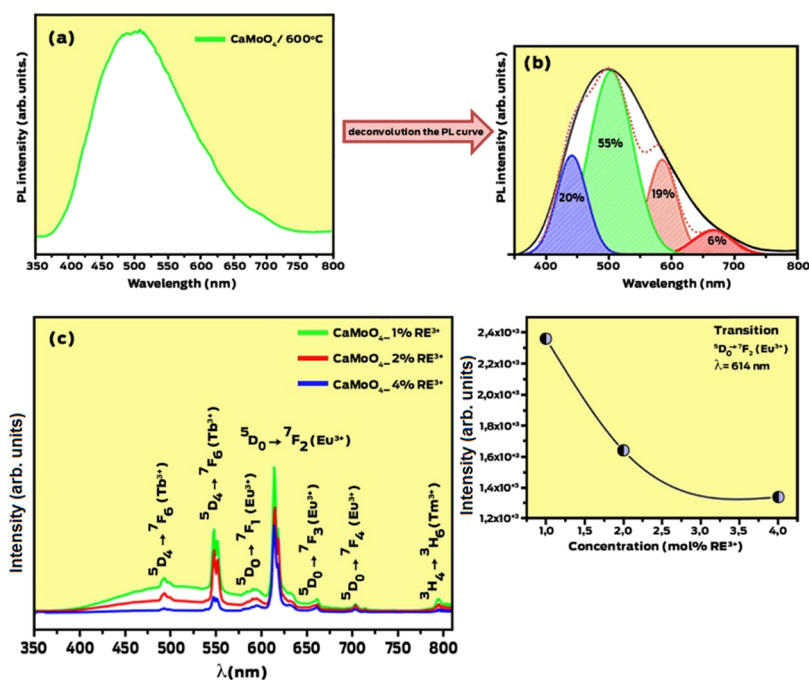


**Figure 6.** Electronic structure analyzed by band structure and total DOS (a)  $\text{Eu}^{3+}$ , (b)  $\text{Tb}^{3+}$ , and (c)  $\text{Tm}^{3+}$  ions for the  $\text{CaMoO}_4:\text{RE}^{3+}$  (12.5%)  $2 \times 2$  supercell.



**Figure 7.** UV–visible absorption spectra for particles: (a)  $\text{CaMoO}_4$ , (b)  $\text{CaMoO}_4:1\%(\text{RE}^{3+})$ , (c)  $\text{CaMoO}_4:2\%(\text{RE}^{3+})$ , and (d)  $\text{CaMoO}_4:4\%(\text{RE}^{3+})$ .





**Figure 8.** PL emission spectrum of (a)  $\text{CaMoO}_4$  at 600 °C, (b) deconvolution of the PL curve at 600 °C, (c)  $(\text{CaMoO}_4:\text{RE}^{3+})$  PL emissions of the powders obtained at 600 °C by the spray pyrolysis method with addition of the  $\text{RE}^{3+}$ . The characteristic color of each  $\text{RE}^{3+}$ ,  $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$  and  $\text{Tm}^{3+}$  is depicted.

The PL spectra of the  $\text{CaMoO}_4:\text{RE}^{3+}$  phosphors (Figure 8c left panel) exhibited characteristics of each dopant–cation ( $\text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ ,  $\text{Tm}^{3+}$ ) emission. Under the excitation of 350 nm, the  $^5\text{D}_4 \rightarrow ^7\text{F}_6$  and  $^5\text{D}_4 \rightarrow ^7\text{F}_5$  transitions occurred at wavelengths of 492 and 548 nm, respectively, which are associated with the emission of  $\text{Tb}^{3+}$ .<sup>117,118</sup> This is related to the CT process from the  $[\text{TbO}_8]$  to the  $[\text{MoO}_4]^{2-}$  tetragonal clusters, which corresponds to the  $4f \rightarrow 5d$  transitions.<sup>119</sup> The emission peak at 800 nm is attributed to  $\text{Tm}^{3+}$ , which is associated with the  $^3\text{H}_4 \rightarrow ^3\text{H}_6$  transition.<sup>115</sup> With respect to  $\text{Eu}^{3+}$ ,  $^5\text{D}_0 \rightarrow ^7\text{F}_j$  ( $j = 1, 2, 3$ , and  $4$ ) transitions occurred at the respective wavelengths of 596, 614, 661, and 704 nm.<sup>117,119</sup> The most intense emission band,  $^5\text{D}_0 \rightarrow ^7\text{F}_2$ , of  $\text{Eu}^{3+}$  is known to have an electric dipole, and this type of transition is highly sensitive to the changes that occur around the  $\text{Eu}^{3+}$  cation. The intensity of different transitions is dependent on the symmetry of the local environment, according to the Judd–Ofelt theory. The analyses of the emission band can provide detailed information on the crystallographic sites occupied by the  $\text{Eu}^{3+}$  cation. The presence of significantly more intense  $^5\text{D}_0 \rightarrow ^7\text{F}_2$  transitions in comparison with the  $^5\text{D}_0 \rightarrow ^7\text{F}_1$  transition indicated a low local symmetry around the  $\text{Eu}^{3+}$  in  $\text{CaMoO}_4$ .

As shown in the right panel of Figure 8c, the emission intensity of the  $\text{RE}^{3+}$  cations reached its maximum at a concentration of 1 mol % of  $\text{RE}^{3+}$  and then decreased with an increase in the concentration. This effect is due to the proximity of the  $\text{RE}^{3+}$  cations within the matrix, which favors the nonradiative transitions. Hence, the low PL signal is enhanced by the quenching effect.<sup>120</sup> The electron interactions that occur between the dopant ions can be of different origins. Their classification is in good agreement with the critical distance (CD) that separates two adjacent  $\text{RE}^{3+}$  cations. The migration of nonradiative energy can occur by two distinct mechanisms: (i) Förster resonance energy transfer (multipole–multipole interaction) and (ii) Dexter mechanism

(exchange interaction). The Förster resonance energy transfer typically occurs at distances of up to 100 Å. The Dexter mechanism (also known as exchange or coalitional energy transfer) is a dynamic quenching mechanism. Dexter energy transfer is a short-range phenomenon ( $\text{CD} = 10$  Å) that decreases with  $e^{-R}$ , and it is dependent on the spatial overlap of donor and quencher molecular orbitals.<sup>120</sup> The critical energy-transfer distance for  $\text{CaMoO}_4:\text{RE}^{3+}$  was estimated using eq 4, as suggested by Blasse,<sup>121</sup> from the parameters of the structure, namely, the unit cell volume ( $V$ ), number of units of the molecular formula per unit cell ( $Z$ ), and quenching concentration ( $X_c$ ).

$$\text{CD} = 2 \left[ \frac{3V}{4X_c \pi Z} \right]^{1/3} \quad (4)$$

With respect to the  $\text{CaMoO}_4:1\% \text{RE}^{3+}$ , the following values were considered:  $Z = 4$ ,  $V = 312.13 \text{ \AA}^3$ , and  $X_c = 0.01$ . The calculated CD was 24.60 Å. In general, the exchange interaction preferably occurs in an energy-transfer process when the value of CD is within the range of 5–10 Å. In the case of values higher than 10 Å, there is no indication of exchange interactions by this mechanism. Consequently, other multipolar electrical interactions are responsible for the quenching effect between two more activating ions ( $\text{RE}^{3+}$ ).

Figure 9 presents the Commission International de l'Éclairage (CIE) chromaticity coordinates for the  $\text{CaMoO}_4:\text{RE}^{3+}$  samples, and Table 2 lists the values for the CIE coordinates, correlated temperature color (CCT), and CRI of the samples. The CIE coordinates were obtained from spectral distributions that define three tristimulus values:  $X$ ,  $Y$ , and  $Z$ . Based on these values, the chromaticity of a certain color can be set using only two coordinates ( $x'$ ,  $y'$ ).

Following the distribution of the PL emissions of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4:\text{RE}^{3+}$  samples, the emitted color of each sample could be characterized. The  $\text{CaMoO}_4$  sample

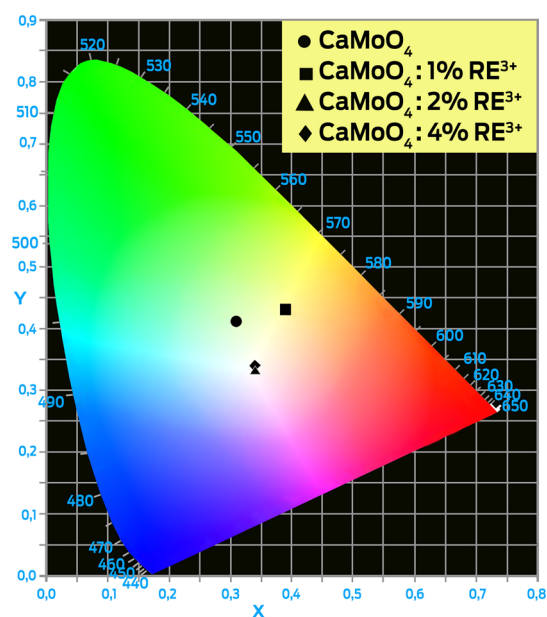


Figure 9. CIE chromaticity diagram of  $\text{CaMoO}_4:\text{RE}^{3+}$ .

released emission with a yellowish green color, whereas the sample doped with 1% of  $\text{RE}^{3+}$  released emission with a greenish yellow color. White emission was observed in the samples doped with 2 and 4% of  $\text{RE}^{3+}$  due to the simultaneous blue, green, and red emissions of the different dopants:  $\text{Tm}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Eu}^{3+}$ , respectively. The region coordinates of the balanced white light of the chromaticity diagram were in the range of ( $x = 0.28\text{--}0.35$ ) and ( $y = 0.30\text{--}0.37$ ).<sup>122</sup> This indicates the presence of a new single phase of a white light emitter. The correlative color temperature was also obtained by the empirical relationship of McCamy.<sup>123</sup> The CCT results confirmed the emission regions, as described by the chromaticity coordinates. All samples reached CRI percentages greater than 90%, which indicates that the samples have a high index of reproduction when used as a light source.

From the surface morphology results based on the FE-SEM measurements, the formation of nonuniform, irregular-sized particles of various dimensions was observed, which could be attributed to the Ostwald ripening of small particles due to the increase in the calcination temperature during the spray pyrolysis process. From the FE-SEM images, the modifications in the morphologies of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4:\text{RE}^{3+}$  microstructures could be identified (Figure 10a–h).

It is believed that the formation of the particle is associated, as the initial solution in which the solvation energy of the  $\text{H}_2\text{O}$  molecule promotes the rapid dissociation of the reactants so that  $\text{Ca}^{2+}$  and  $\text{MoO}_4^{2-}$  ions are rapidly solvated. It is common knowledge that the  $\text{Ca}^{2+}$  cations coordinate to eight oxygen anions, which form a dodecahedral symmetry  $[\text{CaO}_8]$ . The  $\text{Mo}^{6+}$  cations are coordinated with four oxygen anions to form a tetrahedral configuration  $[\text{MoO}_4]$ . Because of the difference

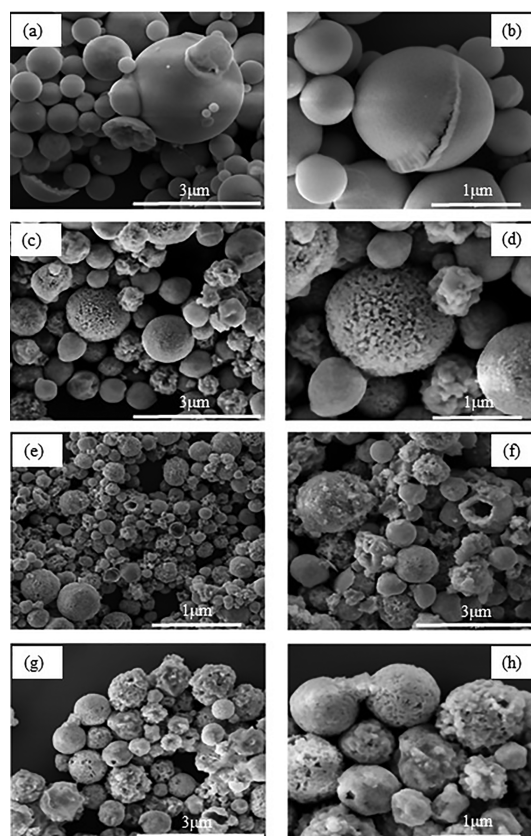


Figure 10. FE-SEM images of  $\text{CaMoO}_4:\text{RE}^{3+}$  (a,b) undoped, (c,d)  $x$ : 1%, (e,f)  $x$ : 2%, and (g,h)  $x$ : 4%.

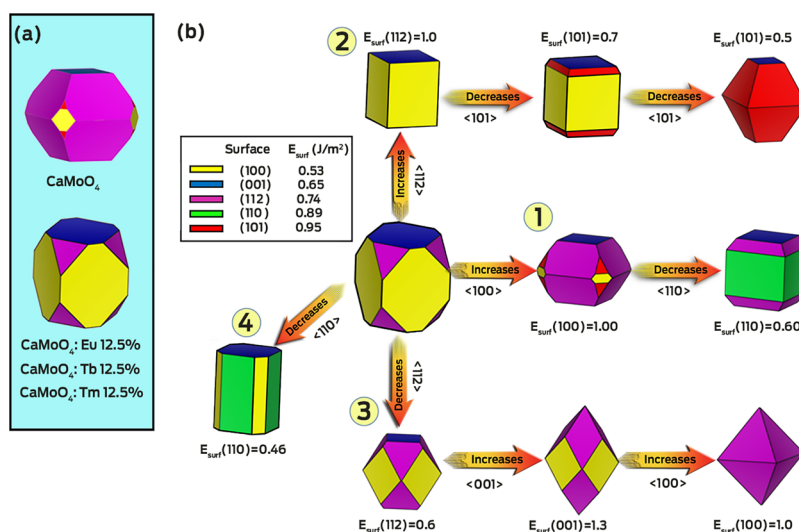
between the electronic densities of the  $[\text{CaO}_8]$  and  $[\text{MoO}_4]$  clusters, there is a strong electrostatic attraction between them. Figure 1 presents a schematic of all steps involved in the growth of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4:\text{RE}^{3+}$  microstructures obtained using the spray pyrolysis method.

For all samples in this study, to which  $\text{RE}^{3+}$  was added (Figure 10c–h), the particles had distorted shapes and exhibited irregular morphologies. Moreover, they were formed by a cluster of polycrystalline particles in which some were dense, and broken particles with holes were observed. All these irregularities in the morphologies were due to the addition of rare earths to the  $\text{CaMoO}_4$  matrix that modify particle surfaces. In addition, the final formation process of the particles using the spray pyrolysis method may promote the emergence of these defects. This was not the case for the particles of pure  $\text{CaMoO}_4$  at 600 °C (Figure 10a,b), which had defect-free surfaces with several cracks due to the rapid formation involved in the pyrolysis spray process.

According to previous calculations, the surface energy of the (001), (112), (110), (101), (100), and (111) surfaces are 0.72, 0.75, 0.93, 1.01, 1.15, and 4.56  $\text{J}/\text{m}^2$ , respectively. Therefore, the order of stability in the pure  $\text{CaMoO}_4$  is as follows: (001)

Table 2. Values of the Chromaticity Coordinates (CIE), CCT, and CRI for  $\text{CaMoO}_4$  and  $\text{CaMoO}_4:\text{RE}^{3+}$  Samples

point symbol	samples	CIE ( $x, y$ )	CCT (K)	CRI	color
●	$\text{CaMoO}_4$	(0.31, 0.41)	6218	91	yellowish green
■	$\text{CaMoO}_4:1\%(\text{RE}/\text{Ca})$	(0.39, 0.43)	4105	93	greenish yellow
▲	$\text{CaMoO}_4:2\%(\text{RE}/\text{Ca})$	(0.34, 0.34)	5171	95	white
◆	$\text{CaMoO}_4:4\%(\text{RE}/\text{Ca})$	(0.34, 0.33)	5095	93	white



**Figure 11.** (a) Theoretical morphology of pure  $\text{CaMoO}_4$  and doped 12.5%  $\text{RE}^{3+}$ . (b) Available morphologies of  $\text{CaMoO}_4$ :12.5%  $\text{Eu}^{3+}$  by tuning the values of  $E_{\text{surf}}$  of the (100), (001), (112), (110), and (101) exposed surfaces.

$< (112) < (110) < (101) < (100)$ .<sup>10</sup> The geometry of the (001), (100), (110), (101), and (112) surfaces of the  $\text{CaMoO}_4$ :12.5%  $\text{Eu}^{3+}$  system are presented in Figure S2 (Supporting Information). Given that the stability of the (111) surface was reduced when compared with the remaining planes in pure  $\text{CaMoO}_4$  because of under-coordination at the  $[\text{MoO}_2]$  cluster,<sup>115</sup> this plane was not evaluated in this study. Moreover, the order of stability of the surfaces was found to be  $(100) < (001) < (112) < (110) < (101)$ . This order was maintained for the three systems at 12.5% with minimal differences between the  $E_{\text{surf}}$  values (see Table S3 in the Supporting Information). The ideal morphologies of the  $\text{CaMoO}_4$ :12.5%  $\text{RE}^{3+}$  microstructures are presented in Figure 11a using Wulff construction, which differs from those of pure  $\text{CaMoO}_4$  due to the large stabilization of the (100) surface. The available morphologies of the  $\text{CaMoO}_4$ : $\text{Eu}^{3+}$  system can be obtained by modifying the relative  $E_{\text{surf}}$  values for each surface, as shown in Figure 11b. In this map, from the ideal morphology and by an increase in the  $E_{\text{surf}}$  value of the (100) surface, a similar morphology (1) to that of pure  $\text{CaMoO}_4$  was achieved; and by an increase in the  $E_{\text{surf}}$  value of the (112) surface, a cubic morphology (2); by a decrease in the  $E_{\text{surf}}$  values for the (112) surface, their presence in the morphology increases, that is (3); and by a decrease in the  $E_{\text{surf}}$  values of the (110) surface, the morphology (4), could be obtained.

In our previous work,<sup>10</sup> as well as in other theoretical and experimental studies<sup>95,96,124,125</sup> the same methodology was proposed to investigate the morphological modulations of different compounds. Therefore, based on this procedure, we can rationalize the morphological transformation imposed by the  $\text{RE}^{3+}$  doping process. Figure 11 displays the main changes caused by rare-earth doping in surface energy values, enabling to predict a plethora of available morphologies for doped  $\text{CaMoO}_4$ .

Finally, it is important to remark that in recent years, a near UV (n-UV) LED chip (350–420 nm) has been extensively investigated, which is combined with blue, green, and red emitting phosphors, to obtain an increased color reproduction index with controllable emitter colors.<sup>126–128</sup> This work presents  $\text{CaMoO}_4$ : $\text{RE}^{3+}$  samples as promising candidates to this end.

## 5. CONCLUSIONS

In summary,  $\text{CaMoO}_4$  and  $\text{CaMoO}_4$ : $x\text{RE}^{3+}$  ( $\text{RE}^{3+} = \text{Eu}^{3+}$ ,  $\text{Tb}^{3+}$ , and  $\text{Tm}^{3+}$ ; and  $x = 1, 2,$  and  $4\%$  mol) compounds with tunable PL emissions and morphologies were successfully synthesized using the spray pyrolysis method. The XRD results revealed that no deleterious phases were present, which confirmed the purity of the samples, and the dopants were then introduced into the matrix. The Raman scattering spectra for the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4$ : $\text{RE}^{3+}$  powders exhibited well-defined peaks for the external and internal modes, and the vibration mode of the scheelite phase was observed in the tetrahedral structures of all samples.

The analysis of the results obtained from first-principles calculation, at DFT level, allows us to rationalize the electronic and optical properties of the as-synthesized samples. In addition, the calculations of the surface energies, based on Wulff construction, explain the morphology transformation of the  $\text{CaMoO}_4$  and  $\text{CaMoO}_4$ : $\text{RE}^{3+}$  microstructures observed in the FE-SEM images. PL measurements point out that the presence of the particular signatures of the  $\text{RE}^{3+}$  cations, that is 615 nm for the  $^5\text{D}_0 \rightarrow ^7\text{F}_2$  transition of  $\text{Eu}^{3+}$ , 545 nm for the  $^5\text{D}_4 \rightarrow ^7\text{F}_5$  transition of  $\text{Tb}^{3+}$ , and 800 nm for the  $^3\text{H}_4 \rightarrow ^3\text{H}_6$  transition of  $\text{Tm}^{3+}$ , reveal a new white light emitter due to the formation of  $\text{CaMoO}_4$ : $\text{RE}^{3+}$  composites. In particular, the chromaticity coordinates of  $\text{CaMoO}_4$ : $\text{RE}^{3+}$  indicated that the samples doped with 2 and 4% of  $\text{RE}^{3+}$  display white emissions, which are promising candidates for visual displays and solid-state lighting. Hence, these materials can demonstrate better performances with certain modifications, thus providing a platform to tune the optical properties.

## ■ ASSOCIATED CONTENT

### 📄 Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcc.9b04123.

Rietveld refined structural parameters; theoretical atomics coordinates; surface energy values; path Brillouin zone and amplified band structure and total DOS projected; surfaces models, chemical analysis of X-ray spectroscopy by dispersive energy; and mapping of the distribution of atoms (PDF)

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## Notes

The authors declare no competing financial interest.

## ACKNOWLEDGMENTS

The authors thank the following Brazilian and Spanish research financing institutions for financial support: National Council for Scientific and Technological Development—CNPq (Processo 303657/2017-0), Sao Paulo Research Foundation—FAPESP (Processo 2013/07296-2), (2016/23891-6), Programa de Pós-Graduação em Ciência e Engenharia de Materiais (PPGCEM-UFRN) and Coordenação de Aperfeiçoamento de Pessoal de Nível Superior—Brasil (CAPES)—Finance Code 001, Universitat Jaume I for project UJIB2016-25, Generalitat Valenciana (PrometeoII/2014/022, ACOMP/2014/270, and ACOMP/2015/1202), Ministerio de Economía y Competitividad, Spain (project CTQ2015-65207-P). M.C.O. acknowledges Generalitat Valenciana for Santiago Grisolia Program 2015/033 and CAPES (2019/88887.319041) for the financial support. We also acknowledge the Servei Informàtica, Universitat Jaume I for a generous allotment of computer time. The authors thank Enio Longo for the support with the scientific illustrations.

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