



# Tailoring WO<sub>3</sub> photoelectrodes with defect-rich MoO<sub>3-x</sub> nanosheets for efficient water splitting reaction

Hugo L. S. Santos · Leticia S. Bezerra<sup>1,2</sup> · Pedro H. C. Camargo<sup>1</sup> · Lucia H. Mascaro<sup>1,2</sup>

Received: 22 January 2025 / Revised: 11 March 2025 / Accepted: 12 March 2025

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## Abstract

Despite its potential for photochemical and photoelectrochemical applications, tungsten trioxide (WO<sub>3</sub>) presents limitations due to its wide bandgap and rapid charge carrier recombination. Here, the photoelectrochemical performance of WO<sub>3</sub> films were enhanced by incorporating defect-rich MoO<sub>3-x</sub> nanosheets. The WO<sub>3</sub> films were produced using a simple polymer-assisted deposition (PAD) method and subsequently modified with defect-rich MoO<sub>3-x</sub> nanosheets, prepared via solvothermal synthesis, by drop-casting. Electronic microscopy revealed that WO<sub>3</sub> exhibited an agglomerated nano-globular structure with several fissures where the MoO<sub>3-x</sub> nanosheets were anchored. In terms of photoelectrochemical performance, the optimal WO<sub>3</sub>/MoO<sub>3-x</sub> film exhibited photocurrent densities of  $1.30 \pm 0.12$  mA cm<sup>-2</sup> and  $3.20 \pm 0.2$  mA cm<sup>-2</sup> under solar simulator and LED 427 nm illumination, respectively, doubling the photocurrent density of bare WO<sub>3</sub>. This enhanced performance was attributed to the formation of a type II heterojunction, which facilitates more efficient charge carrier separation and due to the catalytic enhancement for the oxygen evolution reaction provided by MoO<sub>3-x</sub>.

**Keywords** Tungsten oxide · Defect-rich molybdenum oxide · Photoanode · Photoelectrochemical water splitting

## Introduction

Turning sunlight into clean fuels that can be stored and moved is crucial for a greener future. Photoelectrochemical (PEC) water splitting, a process that utilizes sunlight to split water into hydrogen and oxygen, stands as a core technology for this proposal. Semiconductor materials are essential for efficient PEC water splitting, as they capture sunlight and convert them into electrons and holes. In this context, tungsten trioxide (WO<sub>3</sub>) appears as a promising material for PEC applications. The longer hole diffusion length (~150 nm) compared to  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> (2–4 nm), a better electron mobility (~12 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>) compared to TiO<sub>2</sub> (~0.3 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>), and its valence band energetically located below the energy required for several oxidation reaction, between them the oxygen evolution reaction (OER), enables the application

of WO<sub>3</sub> as a photoanode.[1, 2]. In addition, this material presents excellent stability in acidic aqueous solutions [3].

While WO<sub>3</sub> exhibits favorable optoelectronic properties, its photoelectrocatalytic efficiency towards OER requires significant enhancement. Due to its low absorption coefficient (104–105 cm<sup>-1</sup>), relatively thick films are necessary to ensure sufficient light absorption in this material [4]. However, thicker films can facilitate charge carrier recombination and diminish OER performance. Moreover, peroxide formation on the WO<sub>3</sub> surface during OER can accelerate photocorrosion, leading to reduced stability [3]. Also, the large indirect bandgap energy, ranging from 2.4 to 2.8 eV, depending on the synthesis method, limits its theoretical maximum solar energy efficiency for hydrogen production to approximately 6% and its theoretical photocurrent density to only 5 mA cm<sup>-2</sup> under AM 1.5 G solar illumination [5]. To address these limitations, extensive research has focused on strategies such as morphology engineering, heterojunction formation, and cocatalyst modification.

Films of WO<sub>3</sub> can be synthesized with different morphologies, depending on the synthesis method. In this regard, nanostructured morphologies play an important role in photoelectrocatalytic activity, since this type of structure can improve the diffusion of photogenerated charge carriers,

✉ Lucia H. Mascaro  
lmascaro@ufscar.br

<sup>1</sup> Department of Chemistry, University of Helsinki, A.I. Virtasen Aukio 1, PO Box 55, N0014 Helsinki, Finland

<sup>2</sup> Department of Chemistry, Federal University of Sao Carlos, 13565-905 Sao Carlos, Sao Paulo, Brazil

increase the light absorption by the semiconductor, and increase the number of active sites due to the larger surface area [6, 7]. Polymer-assisted deposition (PAD) is a convenient technique for fabricating porous, adherent  $\text{WO}_3$  films. In this method, a polymer serves as a stabilizing agent, inhibiting particle growth and promoting film adhesion to the substrate [8]. The  $\text{WO}_3$  films produced by PAD often exhibit an irregular, porous morphology consisting of nanospheres, which provides a large surface area [9].

Heterojunction is another effective strategy for improving charge separation efficiency and the photoelectrocatalytic activity of a device. The energy alignment between the conduction band (CB) and valence band (VB) of the semiconductors is the main driving force for electron transfer between the materials, increasing charge flow and reducing recombination.  $\text{WO}_3$  is frequently used to build heterojunctions with other semiconductors to enhance performance, such as  $\text{WO}_3/\text{CuO}$  [10],  $\text{WO}_3/\text{BiVO}_4$  [11],  $\text{WO}_3/\text{g-C}_3\text{N}_4$  [12] and more recently, and  $\text{WO}_3/\text{MoO}_{3-x}$  [13] for photocatalysis application. Recent studies demonstrate that heterojunctions incorporating  $\text{MoO}_{3-x}$  significantly enhance the light absorption and energy conversion efficiency. This enhancement is attributed to the localized surface plasmon resonance (LSPR) exhibited by these particles, which allows them to effectively capture visible light [14–16]. In addition, oxygen vacancies in  $\text{MoO}_{3-x}$  can trap electrons, preventing them from recombining with holes and enhancing the conductivity. This helps to maintain a higher concentration of charge carriers, which can improve the efficiency of photocatalytic processes [13]. Additionally,  $\text{MoO}_{3-x}$  has a narrower bandgap, allowing it to absorb a wider range of light and generate more charge carriers.

Herein, we report a novel  $\text{WO}_3/\text{MoO}_{3-x}$  photoanode material presenting a superior performance for PEC water splitting successfully produced by polymer-assisted and drop-casting method. The optimal  $\text{WO}_3/\text{MoO}_{3-x}$  film exhibited photocurrent densities of  $1.30 \pm 0.12 \text{ mA cm}^{-2}$  and  $3.20 \pm 0.2 \text{ mA cm}^{-2}$  under solar simulator and LED 427 nm illumination, respectively, doubling the photocurrent density of bare  $\text{WO}_3$ . This enhanced performance was attributed to the formation of a type II heterojunction, which facilitates more efficient charge carrier separation and due to the catalytic enhancement for the oxygen evolution reaction provided by  $\text{MoO}_{3-x}$ .

## Experimental

All reagents were of analytical grade and used without any additional purification. Tungsten powder (12  $\mu\text{m}$ , 99.9%) and molybdenum powder (3–7  $\mu\text{m}$ , 99.95%) was purchased from Alfa Aesar, polyethyleneimine (PEI) branched (Mw ~ 25,000 by LS), hydrogen peroxide (30%), and F-doped  $\text{SnO}_2$  (FTO)

substrates with a resistance of  $7 \cdot 10^{-6} \Omega/\text{m}$  were purchased from Sigma-Aldrich.

## Synthesis of $\text{MoO}_{3-x}$

0.192 g (2 mmol) of molybdenum powder was dissolved in 3 mL of  $\text{H}_2\text{O}_2$  30% solution overnight at room temperature producing a yellow solution. Then, 24 ml of ethanol was added to the Mo solution. The ethanolic solution was put on an autoclave and subjected to heat treatment at 150 °C for 12 h with a heating rate of 3 °C/min. The product formed was centrifuged and washed 4 times with water and dried at 60 °C.

## $\text{MoO}_{3-x}/\text{WO}_3$ thin films

0.5 g of tungsten powder was dissolved in 6 mL of  $\text{H}_2\text{O}_2$  18% solution overnight at room temperature. Then, the solution was heated up to 90 °C under reflux for 30 min to decompose the excess peroxide (the color of the solution turns slightly yellow). In another flask, 10 mg of PEI was dissolved in 0.6 mL of water, and then 0.6 mL of the solution of W was added to the PEI solution and homogenized using a vortex mixer producing a yellowish pasty ink. The  $\text{WO}_3$  films were produced by the drop-casting method, prior to the deposition the FTO substrates were cleaned with acetone and isopropanol for 10 min in an ultrasonic cleaner following by  $\text{NH}_4\text{OH}:\text{H}_2\text{O}_2$  1:1 (v/v) solution at 65 °C for 60 min. The films were produced by the drop of 15  $\mu\text{L}$  of the produced ink on 1.0  $\text{cm}^2$  of the substrate. The films were dried at room temperature for 1 h and subjected to heat treatment at 550 °C for 2 h with a heating rate of 5 °C/min. To produce  $\text{MoO}_{3-x}/\text{WO}_3$ , 10  $\mu\text{L}/\text{layer}$  of the  $\text{MoO}_{3-x}$  dispersion 1 mg/mL in isopropanol was dropped under the  $\text{WO}_3$  films and dried at 60 °C for 20 min.

## Physical characterizations

Morphological characterizations were performed by means of scanning electron microscopy (SEM) images acquired in high-resolution field emission using a ZEISS SUPRATM 35 FEG-SEM, a Hitachi S-4800, or on a Philips XL-30 with an acceleration voltage of 5 kV. For the  $\text{WO}_3$  and  $\text{MoO}_{3-x}$  films, transmission electron microscopy (TEM) measurements taken on a Tecnai F20 electron microscope operating at an accelerating voltage of 200 kV were also performed.

Details of the crystal structure were investigated by X-ray diffraction (XRD) using a Rigaku-DMAX2500PC with  $\text{Cu K}\alpha$  as radiation source with a wavelength of 1.5406 Å, working at 40 kV and with beam current fixed at 30 mA. X-ray diffractograms were obtained with a scanning rate of 0.2°  $\text{min}^{-1}$  and a step size of 0.02°.

Elemental analysis was carried out on the benchtop SEM TM4000PLUS II Hitachi equipped with an AZtecOneGO detector of the Silicon Drift type.

The chemical environment and oxidation state of the elements were determined by X-ray photoelectron spectroscopy (XPS) using a Scienta Omicron model ESCA 2SR spectrometer with Al K $\alpha$  (1486.7 eV) as the excitation source. The positions of the peaks were corrected for possible shifts using the adventitious C 1 s peak at 284.8 eV as a reference. For the deconvolution of the high-resolution spectra of the elements, the Shirley function was used as the baseline function and the GL(30) was used for the shape function, which is a convolution of a Lorentzian function with a Gaussian function being 70% Gaussian and 30% Lorentzian. All XPS spectra were treated in the CasaXPS computer program version 2.3.17PR1.1.120.

The optoelectronic characterizations were performed on a NIR-UV-VIS spectrophotometer (Varian, CARY 5G) equipped with a diffuse reflectance module.

### Photoelectrochemical characterization

The electrochemical measurements were carried out in a quartz window cell using an Ag/AgCl/KCl<sub>sat.</sub> as the reference electrode and a platinum wire as counter-electrode. All measurements were carried out in 0.5 mol L<sup>-1</sup> H<sub>2</sub>SO<sub>4</sub> solution at room temperature. A potentiostat/galvanostat (Autolab PGSTAT302N) was used to acquire the electrochemical data. Light emitting diodes (LEDs) of 427 nm and 740 nm (Kessil PR160L) and a solar simulator AM1.5G (LCS-100, Oriel Instruments) were employed as the light sources.

Electrochemical impedance spectroscopy (EIS) was employed to determine the space charge capacitance ( $C_{sc}$ ) of WO<sub>3</sub> and MoO<sub>3-x</sub> materials deposited on FTO.  $C_{sc}^{-2}$  vs. applied potential plots were constructed to extract the flat-band potential ( $E_{fb}$ ) by the Mott-Schottky equation (Eq. 1). All measurements were performed at 1000 Hz applying 10 mV of perturbation amplitude.

$$\frac{1}{C_{sc}^2} = \frac{2}{e\epsilon_r\epsilon_0N_D A^2} \left( E - E_{fb} - \frac{\kappa_B T}{e} \right) \quad (1)$$

The incident photon-to-current efficiency (IPCE) was performed using the monochromator from a Quantum Efficiency/IPCE Measurement Kit (Newport QE-PV-SI) and the measurement was carried out point by point with the wavelength ranging from 350 to 500 nm. All potentials reported in this study were converted to a reversible hydrogen electrode (RHE) using the Eq. 2.

$$E_{RHE} = E_{Ag/AgCl} + E_{Ag/AgCl}^\circ + (0,059 * pH) \quad (2)$$

where  $E_{RHE}$  is the potential measured against the reversible hydrogen electrode (RHE),  $E_{Ag/AgCl}$  is the applied potential measured against the Ag/AgCl reference electrode, and  $E_{Ag/AgCl}^\circ$  is the standard potential of the Ag/AgCl reference electrode.

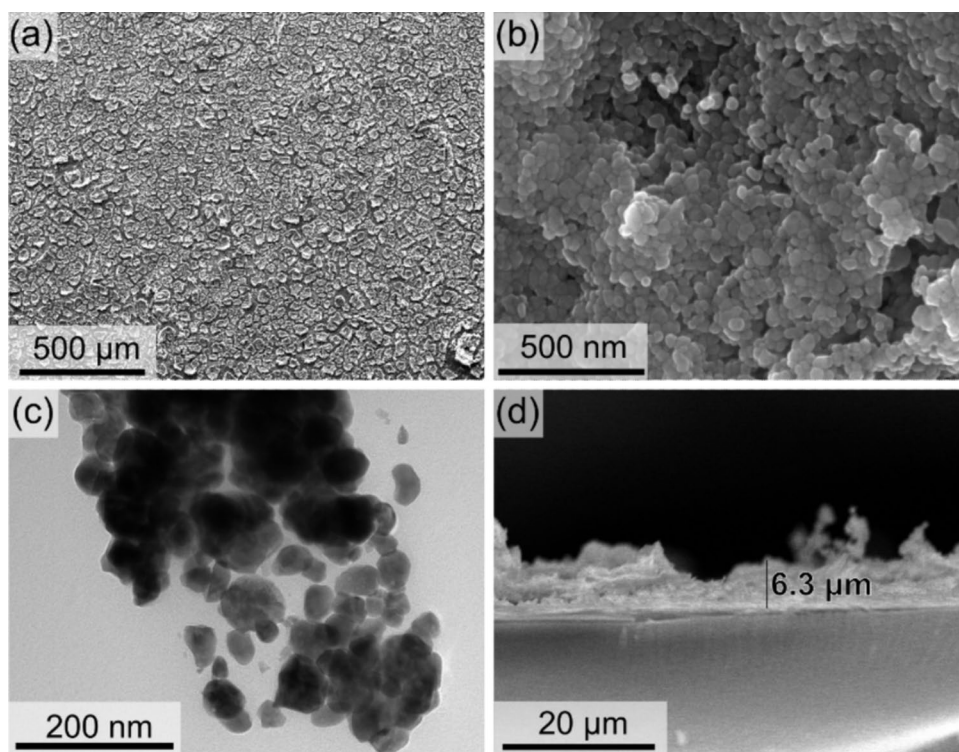
### Results and discussion

The WO<sub>3</sub> films were fabricated using the polymer-assisted deposition (PAD) method. Polyethyleneimine (PEI) worked as a binder and stabilizer for the W species, significantly enhancing film adhesion and surface roughness [9]. The SEM surface image of the WO<sub>3</sub> film at low magnification can be seen in Fig. 1a. The WO<sub>3</sub> film exhibited a highly fractured morphology, with cracks uniformly distributed across its surface. High magnifications SEM image (Fig. 1b) indicated that the WO<sub>3</sub> film was composed of agglomerated nanospheres with an average diameter of  $46 \pm 13$  nm. The nanostructure of WO<sub>3</sub> was also confirmed by TEM as shown in Fig. 1c. The thickness of the WO<sub>3</sub> film was determined from the SEM cross-section as shown in Fig. 1d. Although the exact thickness determination was difficult due to the irregularity of the film, the value of 6.3  $\mu$ m was determined.

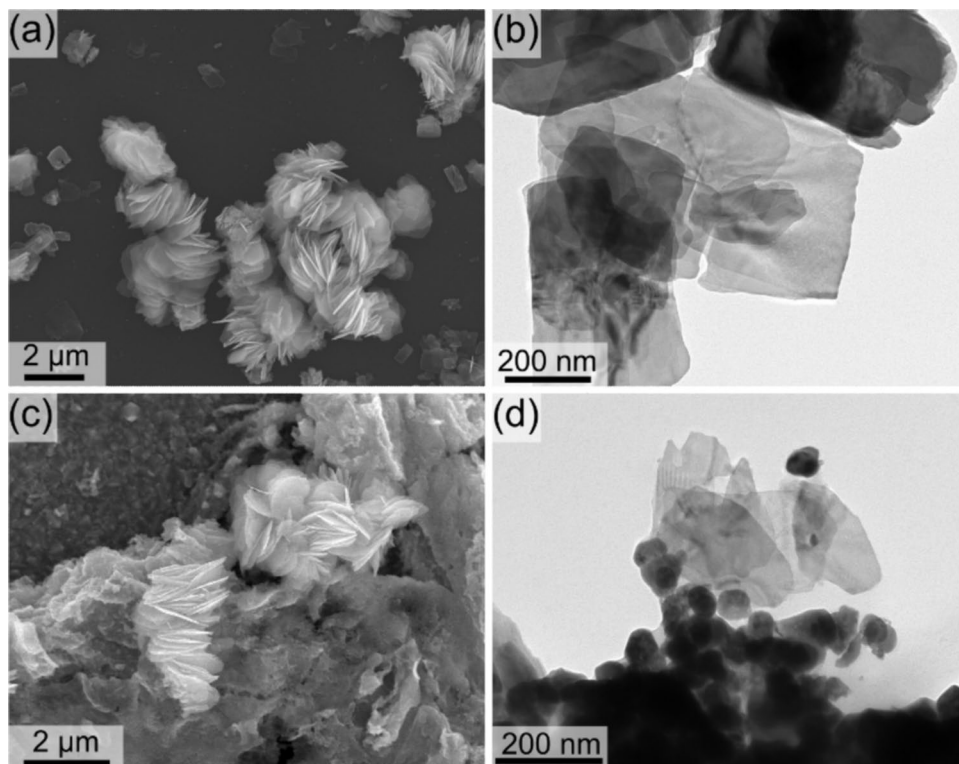
Molybdenum oxide nanostructures (MoO<sub>3-x</sub>) were prepared using a surfactant-free solvothermal synthesis. Molybdenum metal powder was dissolved in hydrogen peroxide, yielding a yellow solution containing complex molybdenum superoxides [17]. The molybdenum superoxide solution was mixed with ethanol and led to hydrothermal treatment in an autoclave reactor. During this process, ethanol partially reduced Mo(VI) species to Mo(V), leading to the formation of non-stoichiometric molybdenum oxide enriched in oxygen vacancies. Morphological characteristics of MoO<sub>3-x</sub> were evaluated by SEM images, as shown in Fig. 2a. The MoO<sub>3-x</sub> particles exhibited a nanosheet morphology with an average thickness of  $25 \pm 7$  nm. These nanosheets were formed through the van der Waals stacking of 2D MoO<sub>3-x</sub> layers [18]. More details of the morphology of MoO<sub>3-x</sub> can be visualized in the TEM image of Fig. 2b.

To elucidate the dispersion of MoO<sub>3-x</sub> nanosheets on the WO<sub>3</sub> film, SEM and TEM imaging was performed on the WO<sub>3</sub>/MoO<sub>3-x</sub> 3L film. As shown in Fig. 2c, the MoO<sub>3-x</sub> nanosheets were predominantly deposited within the cracks of the WO<sub>3</sub> film. Figure 2d provides a magnified view of the interaction between WO<sub>3</sub> and MoO<sub>3-x</sub>. The image reveals MoO<sub>3-x</sub> nanosheets involved by the WO<sub>3</sub> nanospheres, suggesting a strong interfacial contact. The SEM image for the WO<sub>3</sub>/MoO<sub>3-x</sub> 3L film, captured by a backscattered electron (BSE) detector, is presented in Figure S1a. The darker structures, corresponding to the MoO<sub>3-x</sub> nanosheets, are clearly visible within the cracks of the WO<sub>3</sub> film. The elemental

**Fig. 1** **a, b** SEM surface images of  $\text{WO}_3$  films. **c** TEM image for  $\text{WO}_3$ . **d** SEM image of the cross-section of the  $\text{WO}_3$  film



**Fig. 2** SEM (a) and MET (b) images of the  $\text{MoO}_{3-x}$  nanosheets. SEM (c) and MET (d) image of the  $\text{WO}_3/\text{MoO}_{3-x}$  3L film

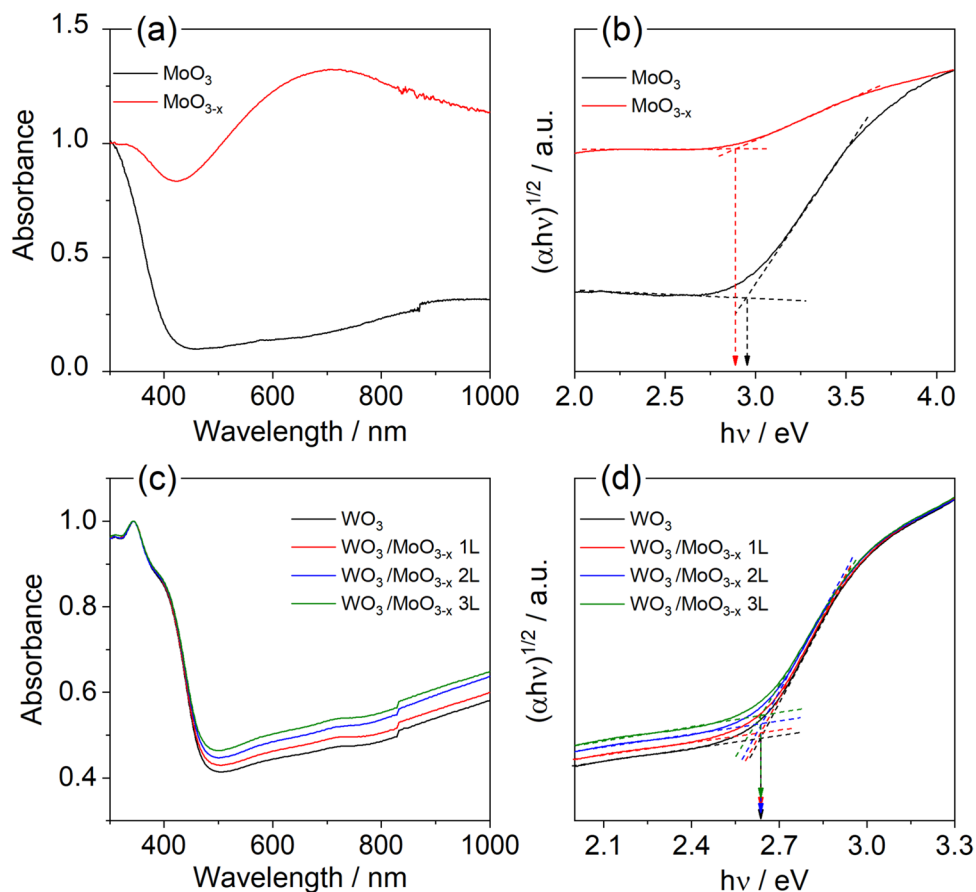


mapping by EDS for W, Mo, and O, shown in Figures S1b, S1c, and S1d, respectively, provide a clear visualization of the elemental distribution.

The optical properties of  $\text{WO}_3$ ,  $\text{MoO}_{3-x}$ , and the  $\text{WO}_3/\text{MoO}_{3-x}$  were investigated through UV–vis spectroscopy. Figure 3a presents the absorbance spectra of  $\text{MoO}_3$  and



**Fig. 3** UV-vis spectra (a) and (c) for the MoO<sub>3</sub>, MoO<sub>3-x</sub>, WO<sub>3</sub>, and WO<sub>3</sub>/MoO<sub>3-x</sub> films. Tauc plots (b) and (d) for the MoO<sub>3</sub>, MoO<sub>3-x</sub>, WO<sub>3</sub>, and WO<sub>3</sub>/MoO<sub>3-x</sub> films



MoO<sub>3-x</sub>. For reference, MoO<sub>3</sub> was synthesized by annealing MoO<sub>3-x</sub> at 350 °C for 2 h in air. MoO<sub>3</sub> exhibits a typical semiconductor absorption profile with an onset around 400 nm, corresponding to the usual interband transitions of a semiconductor. In contrast, MoO<sub>3-x</sub> displays a broad absorption band extending from the near-infrared to the visible region, attributed to localized surface plasmon resonance (LSPR) effects arising from intraband electronic transitions [19, 20]. Additionally, MoO<sub>3-x</sub> nanosheets show a strong absorption band beginning at 405 nm, which can be assigned to interband excitations, similar to MoO<sub>3</sub>. The band gap energy ( $E_g$ ) of MoO<sub>3</sub> and MoO<sub>3-x</sub> was determined from Tauc plots using Eq. 3.[21]

$$(\alpha h\nu)^n \propto (h\nu - E_g) \quad (3)$$

where  $\alpha$  represents the absorption coefficient,  $h$  is Planck's constant,  $\nu$  is the frequency of light, and  $n$  is an exponent that depends on the type of electronic transition. Given the indirect nature of the transitions in WO<sub>3</sub> and MoO<sub>3</sub>,  $n$  is assumed to be 1/2 for both materials [18, 22]. Besides that, to distinguish between direct and indirect band gaps, Tauc plots using the  $n$  values of 1/2 and 2 were performed for the bare WO<sub>3</sub> sample as shown in Figure S2. As observed, the

Tauc plot using the 1/2 (indirect, Figure S2a) presented a very clear linear region which can be used to determine the band gap energy. However, the Tauc for  $n=2$  (direct, Figure S2b) does not have a clear linear relationship between  $(\alpha h\nu)^2$  and  $h\nu$ , confirming this material has an indirect band gap. Band gap ( $E_g$ ) values of  $2.90 \pm 0.03$  eV and  $2.96 \pm 0.05$  eV were obtained for two samples. Although a slight numerical variation is observed, the data indicate no statistically significant difference between the samples, as the observed variation is within the combined experimental error.

Figure 3c presents the absorption spectra of bare WO<sub>3</sub> and WO<sub>3</sub>/MoO<sub>3-x</sub> films. No significant alterations in the interband excitation profile were observed for the WO<sub>3</sub>/MoO<sub>3-x</sub> films compared to pristine WO<sub>3</sub>. However, a noticeable increase in absorption was detected in the visible to near-infrared region (450–800 nm) with an increasing number of MoO<sub>3-x</sub> layers. This enhanced visible light absorption in the WO<sub>3</sub>/MoO<sub>3-x</sub> films can be attributed to the localized surface plasmon resonance (LSPR) effect of MoO<sub>3-x</sub>. The  $E_g$  of bare WO<sub>3</sub> and WO<sub>3</sub>/MoO<sub>3-x</sub> was determined from Tauc plots, as illustrated in Fig. 3d. The  $E_g$  values remained consistent at 2.65 eV for all samples, indicating that the introduction of MoO<sub>3-x</sub> layers did not significantly impact the intrinsic band gap of WO<sub>3</sub>.

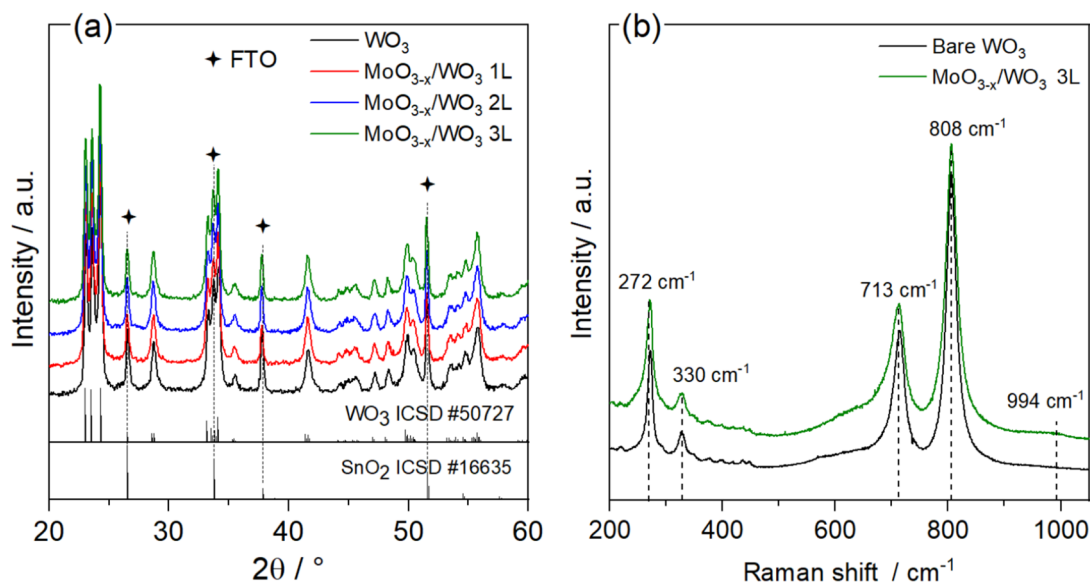
The crystallographic properties of the  $\text{WO}_3$  and  $\text{WO}_3/\text{MoO}_{3-x}$  films were investigated using X-ray diffraction (XRD). The resulting diffractograms are presented in Fig. 4a. The diffraction patterns exhibit peaks corresponding only to the monoclinic  $\text{WO}_3$  phase (ICSD file no. 50727) and the tetragonal  $\text{SnO}_2$  phase originating from the FTO substrate (ICSD no. 16635). The low concentration or amorphous characteristic of  $\text{MoO}_{3-x}$  in the  $\text{WO}_3/\text{MoO}_{3-x}$  films may account for the absence of detectable  $\text{MoO}_3$  peaks in the XRD patterns.

The  $\text{WO}_3$  and  $\text{WO}_3/\text{MoO}_{3-x}$  3C were also characterized by Raman spectroscopy (Fig. 4b). Both films showed four main bands at 272, 330, 713, and 808  $\text{cm}^{-1}$ , which are relative to the vibrational modes of the monoclinic phase of  $\text{WO}_3$ . This corroborates with the XRD diffractograms that exhibited peaks of the monoclinic  $\text{WO}_3$  phase. The main vibrational modes observed, pertain to the stretching of the  $\nu(\text{O}-\text{W}-\text{O})$  and  $\nu(\text{W}-\text{O})$  bonds at 900–600  $\text{cm}^{-1}$  and the bending of the  $\sigma(\text{O}-\text{W}-\text{O})$  bond in the region of 400–200  $\text{cm}^{-1}$  [23, 24]. In addition to the  $\text{WO}_3$  vibrational modes, a broad low-intensity band at 994  $\text{cm}^{-1}$  was identified only in the spectrum of the  $\text{WO}_3/\text{MoO}_{3-x}$  3L film and was assigned to the vibrational mode of stretching of the  $\nu(\text{Mo}=\text{O})$  terminal bond, indicating the presence of  $\text{MoO}_{3-x}$  [25, 26].

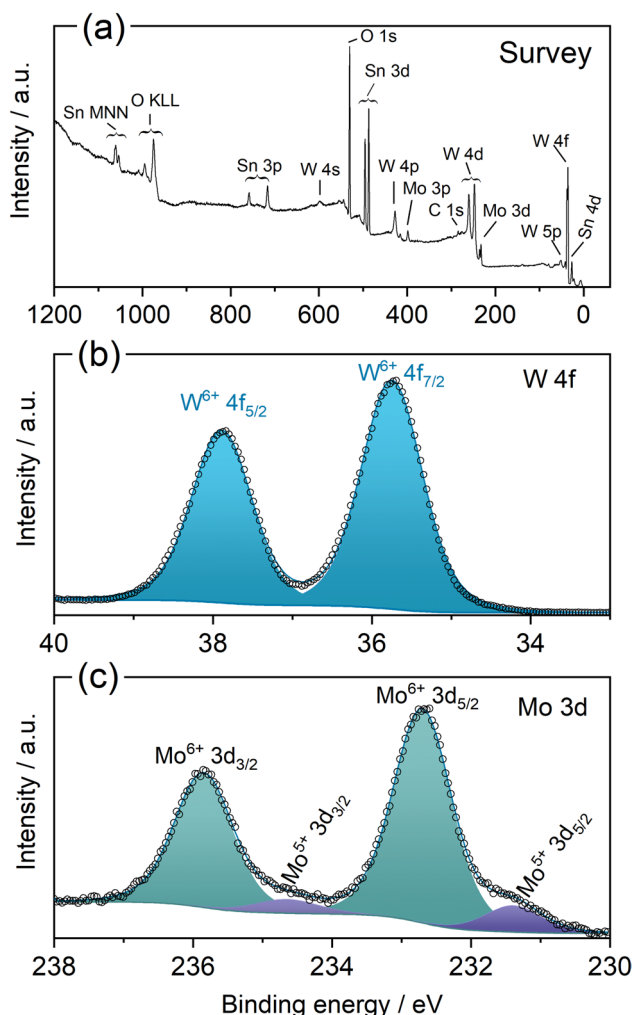
The surface chemical composition and oxidation states of the elements in the  $\text{WO}_3/\text{MoO}_{3-x}$  3L film were determined using X-ray photoelectron spectroscopy (XPS). As shown in the survey spectrum (Fig. 5a), the presence of W, Mo, O, and Sn was confirmed. The inclusion of Sn is attributed to the substrate material. The atomic percentages of W and Mo were calculated to be 85.4% and 14.6%, respectively. The high-resolution XPS spectrum for W 4f (Fig. 5b)

exhibits two prominent peaks at approximately 35.8 eV and 38.0 eV, corresponding to  $\text{W}^{6+} 4f_{7/2}$  and  $\text{W}^{6+} 4f_{5/2}$ , respectively. These findings align with the XRD analysis, which exclusively revealed peaks associated with monoclinic  $\text{WO}_3$  [25, 27]. The high-resolution Mo 4d spectrum (Fig. 5c) was deconvoluted into four distinct peaks at 232.7 eV, 235.9 eV, 231.4 eV, and 234.6 eV, assigned to  $\text{Mo}^{6+} 3d_{5/2}$ ,  $\text{Mo}^{6+} 3d_{3/2}$ ,  $\text{Mo}^{5+} 3d_{5/2}$ , and  $\text{Mo}^{5+} 3d_{3/2}$ , respectively [25, 27]. The atomic percentages of  $\text{Mo}^{6+}$  and  $\text{Mo}^{5+}$  were determined to be 88.2% and 11.8%, respectively. It is important to note that the concentration of  $\text{Mo}^{5+}$  is directly correlated with the concentration of oxygen vacancies (OVs), which is a key factor in the emergence of localized surface plasmon resonance (LSPR) in  $\text{MoO}_{3-x}$  [28, 29].

Linear sweep voltammetry (LSV) and chronopotentiometry were employed to evaluate the photoelectrocatalytic activity of the films for oxygen evolution reaction (OER). All measurements were conducted in a 0.5 M  $\text{H}_2\text{SO}_4$  electrolyte solution. Figure 6a presents the LSV curves for  $\text{WO}_3$  and  $\text{WO}_3/\text{MoO}_{3-x}$  films under simulated sunlight AM 1.5G (100  $\text{mW cm}^{-2}$ ). The voltammograms recorded in the dark (dotted lines) exhibited negligible electrochemical activity, indicating the excellent electrochemical stability of the photoanodes within the studied potential range. In terms of photoelectrocatalytic performance, bare  $\text{WO}_3$  demonstrated the lowest photoelectrochemical activity with  $0.80 \pm 0.05 \text{ mA cm}^{-2}$ . Conversely, the  $\text{MoO}_{3-x}/\text{WO}_3$  3L film exhibited the highest response among all samples, achieving a photocurrent density of  $1.30 \pm 0.12 \text{ mA cm}^{-2}$  at 1.23  $\text{V}_{\text{RHE}}$ , as observed in Table S1. This represents a substantial 60% enhancement compared to bare  $\text{WO}_3$ . LSVs for all electrodes were also conducted under 427 nm LED illumination (50



**Fig. 4** **a** XRD diffractograms for  $\text{WO}_3$  and  $\text{WO}_3/\text{MoO}_{3-x}$  films. **b** Raman spectra for  $\text{WO}_3$  and  $\text{WO}_3/\text{MoO}_{3-x}$  3C films



**Fig. 5** XPS spectra of survey (a), high resolution W 4f (b), and high-resolution Mo 3d for  $\text{WO}_3/\text{MoO}_{3-x}$  3L (c)

$\text{mW cm}^{-2}$ ), as shown in Fig. 6b. The photocurrent generated using the LED was nearly double that observed under simulated solar light. This enhanced photocurrent can be attributed to the superior light absorption properties of  $\text{WO}_3$  and  $\text{MoO}_{3-x}/\text{WO}_3$  under 427 nm LED illumination. The higher energy photons emitted by the LED, which exceed the bandgaps of  $\text{WO}_3$  and  $\text{MoO}_{3-x}$ , result in more efficient excitation of charge carriers. The photoelectrocatalytic activity of all electrodes follows the same trend observed for solar-illuminated experiments, with the  $\text{MoO}_{3-x}/\text{WO}_3$  3L presenting the highest photocurrent of  $3.20 \pm 0.2 \text{ mA cm}^{-2}$  (Table S1).

LSVs were also recorded under 740 nm LED illumination, as shown in Figure S3. The aim of this experiment was to investigate whether the LSPR effect exhibited by  $\text{MoO}_{3-x}$  at 740 nm could enhance the photoelectrocatalytic activity of the  $\text{WO}_3/\text{MoO}_{3-x}$  photoelectrodes. It was hypothesized that the “hot” electrons generated in  $\text{MoO}_{3-x}$  through the LSPR effect could contribute to the OER process. However,

no detectable photocurrent was observed for any of the analyzed films. This lack of photocurrent enhancement may be attributed to the inability of the “hot” electrons generated in  $\text{MoO}_{3-x}$  at 740 nm to be injected into the conduction band of  $\text{WO}_3$  due to energetic barriers correlated to the electronic band position.

The incident photon-to-current conversion efficiency (IPCE) is an important parameter to evaluate the conversion efficiency of the incident photon flux to photocurrent, as a function of the wavelength of the light source. The IPCE can be calculated from Eq. 4 [5, 30].

$$\text{IPCE}\% = \frac{|j(\text{mAcm}^{-2})| \times 1239.8(\text{V} \times \text{nm})}{P_{\text{mono}}(\text{mWcm}^{-2}) \times \lambda(\text{nm})} \quad (4)$$

where  $j$  is the measured photocurrent density, 1239.8 is the product from the multiplication of Planck’s constant ( $h$ ) by the speed of light ( $c$ ),  $P_{\text{mono}}$  is the power of the incident light, and  $\lambda$  is the wavelength of the incident light (nm). The IPCE spectrum, generated using a monochromatic light source, is presented in Fig. 6c. The  $\text{WO}_3/\text{MoO}_{3-x}$  3L film exhibited significantly higher IPCE values, reaching approximately 79% at 410 nm. The most substantial difference in IPCE between the tested films was observed within the 420–430 nm wavelength range, which corresponds to the LED used in the LSV measurements. While 48% of the photons were converted to photocurrent for  $\text{WO}_3/\text{MoO}_{3-x}$ , only 12.5% were converted by the  $\text{WO}_3$  film. This remarkable improvement in performance can be attributed to enhanced interfacial charge transport and more efficient separation of electron–hole pairs [31].

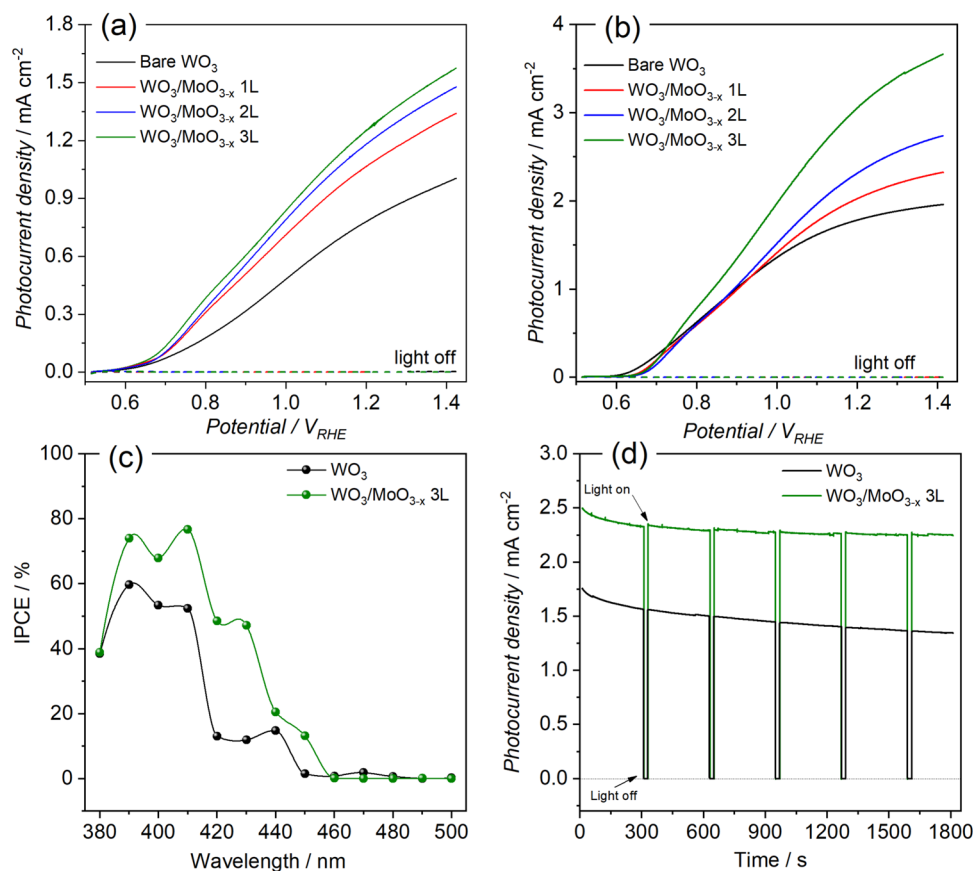
The stability of the pure  $\text{WO}_3$  and  $\text{WO}_3/\text{MoO}_{3-x}$  3L films, a crucial factor in photoelectrocatalysis, was evaluated as shown in Fig. 6d. The photoelectrodes were polarized at 1.1  $\text{V}_{\text{RHE}}$  for 30 min under 427 nm LED illumination. Both photoanodes exhibited excellent stability during the 30-min test, with negligible current observed in the dark. The initial photocurrent retention was determined using Eq. 5.

$$\text{Retention} = \left( \frac{j_f}{j_i} \right) \times 100\% \quad (5)$$

where  $j_i$  and  $j_f$  represent the initial and final photocurrent densities, respectively. The  $\text{WO}_3/\text{MoO}_{3-x}$  3L film exhibited the highest initial photocurrent retention, preserving 89.7% of its photocurrent density after 30 min of reaction. In contrast, pure  $\text{WO}_3$  retained only 76.5%. This indicates that, in addition to enhancing photoelectrocatalytic activity, the incorporation of  $\text{MoO}_{3-x}$  also improves the stability of  $\text{WO}_3$ .

To gain a deeper understanding of the photoelectrochemical performance of  $\text{WO}_3/\text{MoO}_{3-x}$ , the electronic band structures of  $\text{WO}_3$  and  $\text{MoO}_{3-x}$  films were determined. To achieve this, space charge capacitances ( $C_{\text{sc}}$ ) were measured using

**Fig. 6** Linear voltammograms at  $10 \text{ mV s}^{-1}$  of  $\text{WO}_3$  and  $\text{WO}_3/\text{MoO}_{3-x}$  films under illumination of solar simulator (a) and LED 427 nm (b). IPCE efficiencies (c) and stability test (d) at  $1.2 \text{ V}_{\text{RHE}}$  for  $\text{WO}_3$  and  $\text{WO}_3/\text{MoO}_{3-x}$  3L films. Experiments performed in  $0.5 \text{ mol L}^{-1} \text{ H}_2\text{SO}_4$  solution



electrochemical impedance spectroscopy, and  $C_{\text{sc}}$ -voltage plots were plotted. The flat-band potential ( $E_{\text{fb}}$ ) can be correlated to  $C_{\text{sc}}$  using the Mott-Schottky equation (Eq. 1). The  $C_{\text{sc}}$ -voltage plots for  $\text{WO}_3$  and  $\text{MoO}_{3-x}$  films can be seen in Figs. 7a and b, respectively. The straight-line regions were selected based on the observed linearity in the Mott-Schottky plots. These regions correspond to potential ranges where the depletion region capacitance dominates, as evidenced by the consistent linear relationship between  $1/C^2$  and voltage. The selected regions also provided data with minimal noise, ensuring accurate determination of the slope and intercept. Also, to mitigate the discrepancy between the real and geometric surface areas, we fabricated thinner and more compact films of  $\text{WO}_3$  and  $\text{MoO}_{3-x}$  specifically for the M-S analysis. Unlike large nanostructures such as nanowires, the smaller nanospheres of  $\text{WO}_3$  and thin sheets of  $\text{MoO}_{3-x}$  facilitated the creation of compact films with quasi-planar surfaces, minimizing potential errors. The positive slopes in the linear regions confirm the n-type conductivity of these photoelectrodes [18, 32]. The  $E_{\text{fb}}$ , often considered a good approximation of the conduction band edge (CBE) for n-type semiconductors, was determined to be  $0.43 \text{ V}_{\text{RHE}}$  for  $\text{WO}_3$ , which is consistent with previously reported values [11]. For  $\text{MoO}_{3-x}$ , the  $E_{\text{fb}}$  value was measured as  $0.15 \text{ V}_{\text{RHE}}$ , also aligning with literature reports [33]. The valence band edge (VBE) was determined using the

equation  $\text{VBE} = E_{\text{fb}} - E_g$ . Table 1 summarizes the values of  $E_g$ ,  $E_{\text{fb}}$ , CBE, and VBE for the investigated materials.

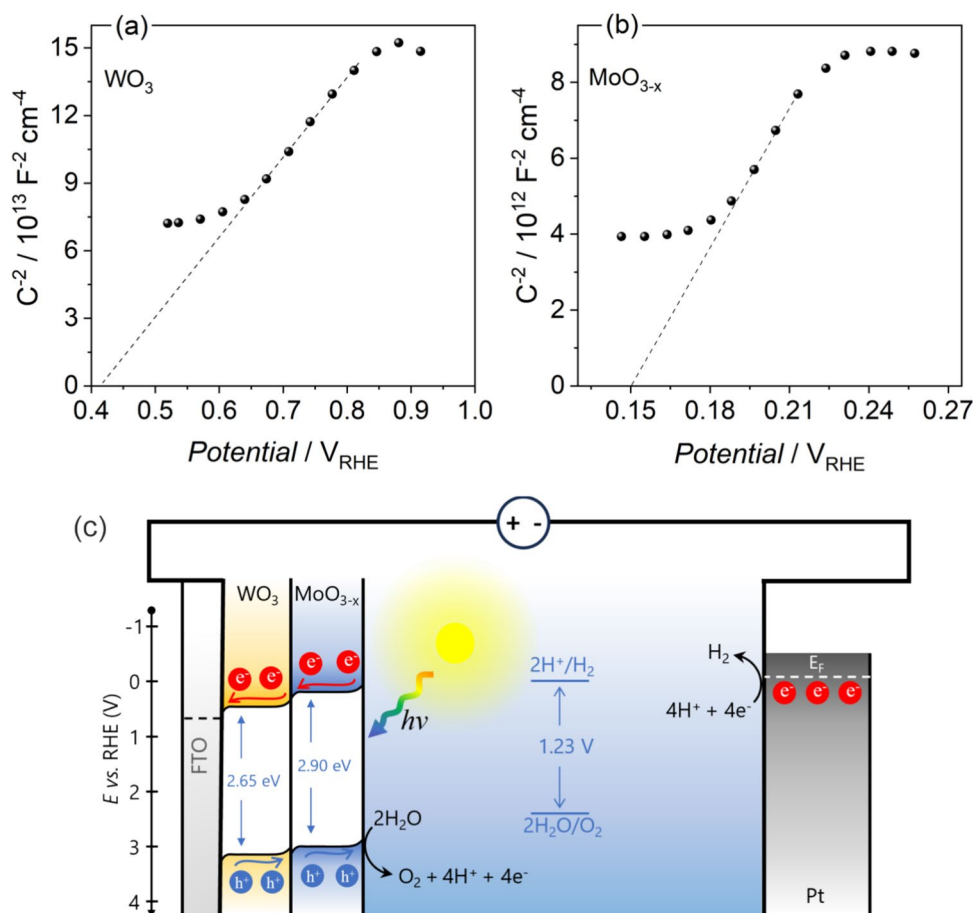
An energy band diagram for the  $\text{WO}_3/\text{MoO}_{3-x}$  heterojunction was constructed based on the calculated CBE and VBE values, as depicted in Fig. 7c. The relative positions of these energy levels indicate the formation of a type II n-n heterojunction. This configuration facilitates efficient charge separation due to the staggered alignment of the conduction and valence bands. The resulting potential gradient further enhances the transport of photogenerated charge carriers from the  $\text{WO}_3$  to the  $\text{MoO}_{3-x}$ , where they can participate in redox reactions. Furthermore,  $\text{MoO}_{3-x}$  may have acted also as a co-catalyst improving the overall photoelectrocatalytic performance of the  $\text{WO}_3/\text{MoO}_{3-x}$  heterojunction. The synergistic combination of these factors, including efficient charge separation, enhanced charge transport, and accelerated reaction kinetics, contributes to the superior photoelectrocatalytic performance of the  $\text{WO}_3/\text{MoO}_{3-x}$  compared to pristine  $\text{WO}_3$ .

## Conclusion

An efficient  $\text{WO}_3/\text{MoO}_{3-x}$  heterojunction was successfully fabricated using straightforward techniques such as polymer-assisted deposition, solvothermal synthesis, and



**Fig. 7** Mott-Schottky plots at 1000 Hz for WO<sub>3</sub> e (a) MoO<sub>3-x</sub> in 0.5 mol L<sup>-1</sup> H<sub>2</sub>SO<sub>4</sub> (b). c Illustration of the energy band diagram for the WO<sub>3</sub>/MoO<sub>3-x</sub> heterojunction



**Table 1** Optical  $E_g$ ,  $E_{fb}$ , CBE, and VBE values for WO<sub>3</sub> and MoO<sub>3-x</sub> films

Sample	$E_g$ (eV)	$E_{fb}$ (V <sub>RHE</sub> )	CBE (V <sub>RHE</sub> )	VBE (V <sub>RHE</sub> )
WO <sub>3</sub>	2.65	0.43	0.43	3.08
MoO <sub>3-x</sub>	2.90	0.15	0.15	3.05

drop-casting. By optimizing the number of MoO<sub>3-x</sub> layers to three, the resulting WO<sub>3</sub>/MoO<sub>3-x</sub> film demonstrated significant enhancement in photoelectrochemical performance, achieving photocurrent densities of 1.30 mA cm<sup>-2</sup> and 3.18 mA cm<sup>-2</sup> under simulated sunlight and 427 nm LED illumination, respectively. This substantial improvement, more than doubling the photocurrent density of pristine WO<sub>3</sub>, can be attributed to the formation of a type II heterojunction. This configuration facilitates efficient charge carrier separation, preventing recombination losses and promoting charge transport to the surface. Additionally, the incorporation of MoO<sub>3-x</sub> may have accelerated the kinetics of the rate-determining step, leading to enhanced overall photoelectrochemical performance. These findings highlight the potential of the WO<sub>3</sub>/MoO<sub>3-x</sub> heterojunction

as a promising candidate for efficient solar energy conversion and water splitting applications.

**Supplementary Information** The online version contains supplementary material available at <https://doi.org/10.1007/s10008-025-06277-2>.

**Acknowledgements** The authors thank the São Paulo Research Foundation (FAPESP) for the financial assistance to the project and the fellowship granted (#2021/08184-0, #2019/26860-2, #2017/11986-5, and #2013/07296-2); the Conselho Nacional de Desenvolvimento Científico e Tecnológico (CNPq) (#406156/2022-0, #311769/2022-5); FINEP-MCTI 01.22.0179.00 and Shell and the strategic importance of the support given by ANP (Brazil's National Oil, Natural Gas and Biofuels Agency) through the R&D levy regulation.

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